

ENERGY EQUALITY IN COMPRESSIBLE FLUIDS WITH PHYSICAL BOUNDARIES

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ABSTRACT. We study the energy balance for the weak solutions of the three-dimensional compressible Navier–Stokes equations in a bounded domain. By constructing a global mollification combined with an independent boundary cut-off, and then taking a double limit to prove the convergence of the resolved energy, we establish an L^p - L^q regularity condition on the velocity field for the energy equality to hold, provided that the density is bounded and satisfies $\sqrt{\rho} \in L_t^\infty H_x^1$. As a result of our new approach, we can avoid assuming additional regularity of the velocity near the boundary in order to deal with the boundary production due to the diffusion terms.

1. INTRODUCTION

In fluid mechanics, compressible fluids play an important role in many fields of applications, including astrophysics (star-formation, interstellar/intergalactic medium), engineering (supersonic aircraft, gas turbines, combustion engines), and so on. In this paper we consider the following three-dimensional Navier–Stokes equations of isentropic compressible flows, consisting of the conservation of mass and momentum,

$$\begin{cases} \rho_t + \operatorname{div}(\rho u) = 0, \\ (\rho u)_t + \operatorname{div}(\rho u \otimes u) + \nabla P(\rho) - \mu \Delta u - (\mu + \lambda) \nabla \operatorname{div} u = 0, \end{cases} \quad (1.1)$$

where $\rho \geq 0$ is the density of the flow, $u \in \mathbb{R}^3$ is the velocity, and $P(\rho) = \rho^\gamma$ is the pressure with $\gamma > 1$. The viscosity constants include the shear viscosity $\mu > 0$ and the bulk viscosity λ satisfying $\lambda + \frac{2}{3}\mu \geq 0$.

We are particularly interested in the behavior of the compressible flows confined within solid walls. Such flows are ubiquitous in nature as well as in applications. Mathematically we consider the above system (1.1) in an open bounded domain $\Omega \subset \mathbb{R}^3$ and pose the usual no-slip boundary condition

$$u = 0 \quad \text{on } \partial\Omega. \quad (1.2)$$

Finally we complement (1.1) with the initial condition

$$\rho(x, 0) = \rho_0(x), \quad (\rho u)(x, 0) = (\rho_0 u_0)(x), \quad x \in \Omega, \quad (1.3)$$

where we define $u_0 = 0$ on the sets $\{x \in \Omega : \rho_0 = 0\}$.

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System (1.1)–(1.3) possesses an energy balance law that holds at least formally for strong solutions:

$$\begin{aligned} \int_{\Omega} \left(\frac{1}{2} \rho |u|^2 + \frac{\rho^\gamma}{\gamma - 1} \right) dx + \int_0^t \int_{\Omega} (\mu |\nabla u|^2 + (\mu + \lambda) (\operatorname{div} u)^2) dx ds \\ = \int_{\Omega} \left(\frac{1}{2} \rho_0 |u_0|^2 + \frac{\rho_0^\gamma}{\gamma - 1} \right) dx. \end{aligned} \quad (1.4)$$

On the other hand, from the classical results of Lions [28] and Feireisl [18, 19], this system also allows for solutions with less regularity, namely the weak solutions (see below), which only satisfy an energy inequality.

Definition 1.1. For a given $T > 0$, we call (ρ, u) a weak solution on $[0, T]$ to (1.1)–(1.3) if

- The problem (1.1)–(1.3) holds in $\mathcal{D}'([0, T] \times \Omega)$ and

$$\rho^\gamma, \rho |u|^2 \in L^\infty(0, T; L^1(\Omega)), \quad u \in L^2(0, T; H_0^1(\Omega)). \quad (1.5)$$

- (ρ, u) is a renormalized solutions of (1.1)₁ in the sense of [9].
- The energy inequality holds

$$\begin{aligned} \int_{\Omega} \left(\frac{1}{2} \rho |u|^2 + \frac{\rho^\gamma}{\gamma - 1} \right) dx + \int_0^t \int_{\Omega} (\mu |\nabla u|^2 + (\mu + \lambda) (\operatorname{div} u)^2) dx ds \\ \leq \int_{\Omega} \left(\frac{1}{2} \rho_0 |u_0|^2 + \frac{\rho_0^\gamma}{\gamma - 1} \right) dx. \end{aligned} \quad (1.6)$$

The lack of the exact equality in (1.6) is reminiscent of the energy inequality of the Leray–Hopf solution to the incompressible Navier–Stokes equations, which still remains open up to date. One of the main difficulties in establishing the energy equality in the absence of the boundary (i.e. $\Omega = \mathbb{R}^3$ or \mathbb{T}^3) lies in the fact that the regularized velocity field and density may generate a non-vanishing energy flux due to the nonlinear coupling. For incompressible flows with constant density, J. L. Lions [26] proved that energy equality holds for $u \in L_{t,x}^4$. This was reproduced by Ladyženskaja et al. [22] in the general context of parabolic equations. In [32] Serrin gave a dimension-dependent condition $u \in L_t^p L_x^q$ for $\frac{2}{p} + \frac{n}{q} \leq 1$, $q > n$, where n is the space dimension. Later Shinbrot in [33] removed the dimensional dependence and improved the conditions to $\frac{2}{p} + \frac{2}{q} \leq 1$, $q \geq 4$. An alternative proof of Shinbrot’s result can be found in [36]. New types of conditions have been obtained recently, including Besov-type regularity conditions [6, 13], weak-in-time with optimal Onsager spatial regularity conditions [7], new $L_t^p L_x^q$ conditions in combination with low dimensionality of the singular set [25], to name a few. For inhomogeneous incompressible flows Leslie-Shvydokoy [24] proved the energy equality in Besov spaces. Concerning compressible fluids, the theoretical study is more recent. Using the approach of [8, 14, 17] in the framework of Onsager’s theory [31], Drivas–Eyink [11, 12] derived necessary conditions for dissipative anomalies of kinetic energy in turbulent solutions of the compressible Euler equations. Feireisl et al. [20] gave sufficient Besov regularity conditions on the weak solutions for energy conservation of the compressible Euler system, excluding the case of vacuum. Regularity conditions for energy conservation which allow the presence of vacuum in the compressible Euler flow were provided by Akramov et al. [1]. For the energy equality of the compressible Navier–Stokes (1.1), Yu [35] proved that (1.4) holds

true if the velocity variable u satisfies an $L_t^p L_x^q$ condition while the density ρ is bounded and $\sqrt{\rho} \in L^\infty(0, T; H^1)$.

The presence of solid boundaries makes the dissipative mechanism more complex. The near-wall behavior of the solution may differ substantially from the behavior in the interior bulk. Therefore mathematically, the added challenge comes from controlling the regularity of the solutions near the boundary in order to pass from local to global energy balance. The first result addressing the Onsager's theory for wall-bounded flows is due to Bardos–Titi [3] in the context of the incompressible Euler equations under the assumption of a global regularity on the velocity. Such a result was further refined by Bardos–Titi–Wiedemann [4] and Drivas–Nguyen [10] where a weaker assumption is used that is consistent with the formation of the boundary layer in the vanishing viscosity limit. In line with the method of [4], Akramov et al. [1] were able to treat the case of compressible Euler flows confined in a bounded domain. In [2], Bardos et al. managed to extend and prove the Onsager conjecture for a class of conservation laws that admit a generalized entropy. The idea of [3] was also exploited by Yu [37] for the incompressible Navier–Stokes equations in a bounded domain, obtaining the same Shinbrot type interior regularity criterion, with an additional Besov regularity on the velocity to handle the boundary effects coming from the diffusion term.

The basic strategy used in [1, 3, 4, 10, 37] is localization. Specifically, an additional cut-off function was introduced that separates the boundary part from the interior domain. The distance h from the support of the boundary cut-off to the boundary is chosen to be large enough compared with the scale ε for the mollification, leaving enough space to mollify the interior velocity. This way the interior regularity criterion can be achieved following the classical commutator estimates in the spirit of [8, 14, 17]. To obtain the global energy balance, one needs to patch the interior estimates with the estimates on the boundary layer. This is done by carefully examining the scale-transfer terms in the bulk and at the boundary. For the incompressible Euler equations, to ensure that the inertial boundary production vanishes in the double limit $\varepsilon, h \rightarrow 0$, one needs to assume continuity of the normal component of the energy flux near the boundary [4], which is essentially equivalent to assuming continuity of the near-wall normal velocity [10] due to the non-penetration boundary condition. The case for the incompressible Navier–Stokes is slightly more delicate. The boundary production includes an additional contribution coming from the diffusion term, which involves the information about the velocity gradient near the boundary. However, such information cannot be inferred from the no-slip boundary condition, and this is the reason why in [37] an extra Besov regularity on the velocity is assumed.

1.1. Methodology. The goal of this paper is to understand the relation between the energy equality and the regularity of the solutions in the appearance of the boundary. For this purpose we shall introduce a new approach different from [3] and apply it to the compressible Navier–Stokes system (1.1)–(1.3). As a result of our new approach, we can avoid assuming additional regularity of the velocity near the boundary as in [37] in order to deal with the boundary production due to the diffusion terms. To the best of our knowledge, our paper appears to be the first work addressing the energy balance of flows both in the compressible regime and in a bounded domain.

In the paragraphs below we briefly describe the ideas of our method.

Global mollification. The approach we propose in this paper is “global” in the sense that we do not shrink the domain Ω to create space for the mollification. Instead, the mollified functions are defined globally in the whole domain Ω . Roughly speaking, the interior mollification will be the same as in the general localization approach. However, for the boundary part, when $\partial\Omega$ is reasonably smooth, we introduce a local variable shift toward the interior of Ω and then perform the usual mollification. Finally we obtain a global approximation by gluing together the boundary and interior parts using a partition of unity. The details are given in Section 2.1. We want to point out that such an approximation is in the spirit of the one discussed in [16, Section 5.3].

The regularization of the momentum equation in (1.1) can be done the same way: performing local mollifications, and then summing them up according to the partition of unity.

Test functions. The global approximation avoids cutting out the boundary information, at the price that the mollified velocity field fails to vanish on the boundary. Therefore one still needs to introduce a boundary cut-off function supported δ -distance away the boundary (cf. (3.4)), and multiply it to the mollified velocity to construct the test function. The difference, compared with [3, 4, 10, 37], is that the mollification scale ε and boundary cut-off scale δ are completely independent. This leaves much freedom for the choice of δ and could be useful for other applications, for instance, the study of anomalous dissipation in the vanishing viscosity limit, which will be addressed in a forthcoming paper [5].

Boundary production due to diffusion. Similarly to [37], the inertial boundary production includes terms that involve the gradient of the velocity field ∇u which comes from the diffusion terms. As explained earlier, it is hard to control such terms directly due to the lack of boundary condition on ∇u . Here we will first pass the limit as $\varepsilon \rightarrow 0$, leaving δ fixed, so that we recover the full velocity in the resulting approximated energy equality (3.27). This allows us to employ the classical Hardy type inequality (cf. Lemma 2.3) to annihilate the boundary contribution from the diffusion terms. Note that the crucial ingredient in this argument is the fact that ε and δ are independent.

Commutator estimates. In proving energy conservation/equality, the commutator estimates are required for treating the nonlinear terms. Compared with incompressible homogeneous equations, a notable difference in compressible (or inhomogeneous) equations is that the momentum equation contains a time derivative of a nonlinear term ρu , and hence it needs a commutator estimate in time. We follow the ideas in [35] in order to allow for vacuum states, with slight modifications to work in the Sobolev spaces; see Corollary 2.1.

1.2. Main results. Our energy equality criterion for the compressible Navier–Stokes equations (1.1)–(1.3) is

Theorem 1.1. *Let Ω be an open, bounded domain with C^1 boundary $\partial\Omega$, and (ρ, u) be a weak solution in Definition 1.1. Assume that*

$$0 \leq \rho \leq \bar{\rho} < \infty, \quad \nabla \sqrt{\bar{\rho}} \in L^\infty(0, T; L^2(\Omega)). \quad (1.7)$$

If

$$u \in L^p(0, T; L^q(\Omega)), \quad p \geq 4, q \geq 6, \quad (1.8)$$

and moreover,

$$u_0 \in L^{q_0}(\Omega), \quad q_0 > 3. \quad (1.9)$$

Then (1.4) holds for any $t \in [0, T]$.

A few remarks are in order as follows.

Remark 1.1. Condition (1.8) can be improved in the absence of the vacuum states to

$$u \in L^p(0, T; L^q(\Omega)), \quad \text{with} \quad \frac{2}{p} + \frac{3}{2q} \leq \frac{3}{4}, \quad q \geq 6. \quad (1.10)$$

In fact, it follows from (1.5) and (1.7) that $u \in L^\infty(0, T; L^2(\Omega))$, which implies by interpolation

$$\|u\|_{L^4(0, T; L^6)} \leq \|u\|_{L^\infty(0, T; L^2)}^{\frac{(q-6)}{3(q-2)}} \|u\|_{L^{\frac{8q}{3(q-2)}}(0, T; L^q)}^{\frac{2q}{3(q-2)}} \leq C \|u\|_{L^p(0, T; L^q)},$$

as long as $\frac{8q}{3(q-2)} \leq p$. This is condition (1.10).

Remark 1.2. We will apply the same idea to treat the Leray–Hopf solution of the incompressible Navier–Stokes equations in a bounded domain in the Appendix. We are able to obtain an analogous regularity criterion as in the periodic case, with an additional condition on the control of the pressure on the boundary. This removes the extra Besov regularity assumption on the velocity as in [37].

Remark 1.3. The regularity assumption (1.7) on the density is critical for making commutator estimates work, but it is not optimal. Alternatively, it can be relaxed at the expense of imposing extra time regularity on velocity field. This is similar to, for e.g., [12, 20].

The rest of this paper is organized as follows: In Section 2 we construct the global mollification, prove the commutator estimates in Sobolev spaces, and recall a classical Hardy-type inequality. In Section 3 we give the proof for the main theorem. Finally in the Appendix we apply our method to the incompressible Navier–Stokes equations in a bounded domain and give sufficient regularity conditions for the energy equality.

2. PRELIMINARIES

2.1. Global approximation in Ω . If $f \in L^p(0, T; W^{1,p}(\Omega))$, the following local approximation is well known

$$f^\varepsilon \rightarrow f \quad \text{in} \quad L_{loc}^p(0, T; W_{loc}^{1,p}(\Omega)), \quad \forall p \in [1, \infty), \quad (2.1)$$

where

$$f^\varepsilon(x, t) = \int_0^T \int_\Omega f(y, s) \eta_\varepsilon(x - y, t - s) dy ds, \quad \eta_\varepsilon(x, t) = \frac{1}{\varepsilon^4} \eta\left(\frac{x}{\varepsilon}, \frac{t}{\varepsilon}\right), \quad (2.2)$$

with $\eta(x, t)$ being the standard mollifier supported in a unit ball.

For the purpose of this paper, we adopt some ideas in [16, Section 5.3] and build a global approximation in $L_{loc}^p(0, T; W^{1,p}(\Omega))$.

(1) Since $\partial\Omega \in C^1$, for a fixed $x_1 \in \partial\Omega$, there exist some $r_1 > 0$ and a C^1 function $h : \mathbb{R}^2 \rightarrow \mathbb{R}$ such that, upon relabelling the coordinate axis if necessary, we have

$$\Omega \cap B(x_1, r_1) = \{x \in B(x_1, r_1) : x_3 > h(x_1, x_2)\},$$

where $B(x, r)$ is an open ball which centers in x with radius r .

Let $V_1 = \Omega \cap B(x_1, \frac{r_1}{2})$. For a small $\varepsilon < \frac{r_1}{8}$, we define the shifted point

$$x^\varepsilon := x - \varepsilon \vec{n}(x_1), \quad (2.3)$$

then it is obvious that

$$B(x^\varepsilon, \varepsilon) \subset \Omega \cap B(x_1, r_1), \quad \text{for all } x \in V_1,$$

where $\vec{n}(x_1)$ is the unit outward normal vector of $\partial\Omega$ at x_1 (see Fig. 1 below).

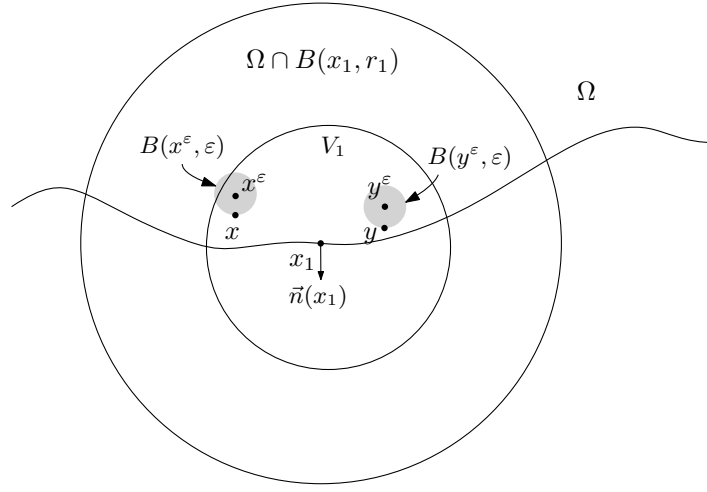


FIGURE 1. The local variable shift defined in V_1 .

Define the shifted function

$$\tilde{f}(x, t) = f(x^\varepsilon, t), \quad x \in V_1. \quad (2.4)$$

Then there is room to mollify $\tilde{f}(x, t)$ like (2.2), that is,

$$\begin{aligned} \tilde{f}_1^\varepsilon(x, t) &= \int_0^T \int_{\mathcal{V}_1} \tilde{f}(y, s) \eta_\varepsilon(x - y, t - s) dy ds \\ &= \int_0^T \int_{\mathcal{V}_1 - \varepsilon \vec{n}(x_1)} f(y, s) \eta_\varepsilon(x^\varepsilon - y, t - s) dy ds, \end{aligned} \quad (2.5)$$

for every $(x, t) \in V_1 \times (\varepsilon, T - \varepsilon)$, and \mathcal{V}_1 can be simply taken to be $B(x_1, \frac{r_1}{2} + 2\varepsilon) \cap \Omega$.

We claim that

$$\lim_{\varepsilon \rightarrow 0} \|\tilde{f}_1^\varepsilon - f\|_{L_{loc}^p(0, T; W^{1, p}(V_1))} = 0. \quad (2.6)$$

To confirm this, for any multi-index α satisfying $|\alpha| \leq 1$,

$$\|\partial_x^\alpha(\tilde{f}_1^\varepsilon - f)\|_{L_{loc}^p([0, T] \times V_1)} \leq \|\partial_x^\alpha(\tilde{f}_1^\varepsilon - \tilde{f})\|_{L_{loc}^p([0, T] \times V_1)} + \|\partial_x^\alpha(\tilde{f} - f)\|_{L_{loc}^p([0, T] \times V_1)}.$$

The second term on the right-hand side of the above goes to zero as $\varepsilon \rightarrow 0$ because the translation is continuous in L^p , and the first term also vanishes as ε goes to zero due to (2.1).

(2) Since $\partial\Omega$ is compact, we find finitely many points $x_i \in \partial\Omega$, radii $r_i > 0$, corresponding sets $V_i = \Omega \cap B(x_i, \frac{r_i}{2})$ and functions $\tilde{f}_i^\varepsilon \in C^\infty(\overline{V}_i)$ ($i = 1, \dots, k$) such that $\partial\Omega \subset \bigcup_{i=1}^k \overline{V}_i$. By (2.6), it follows that for a given $\delta > 0$,

$$\|\tilde{f}_i^\varepsilon - f\|_{L_{loc}^p(0,T;W^{1,p}(V_i))} < \frac{\delta}{k+1}, \quad i \in \{1, \dots, k\}, \quad (2.7)$$

as long as ε is taken small. Take also an open set $V_0 \subset\subset \Omega$ to satisfy

$$\Omega \subset \bigcup_{i=0}^k V_i \quad \text{and} \quad \|f^\varepsilon - f\|_{L_{loc}^p(0,T;W^{1,p}(V_0))} < \frac{\delta}{k+1}. \quad (2.8)$$

We have the following

Proposition 2.1 (Global mollification in Ω). *Let $\{\xi_i\}_{i=0}^k$ be a smooth partition of unity subordinate to V_i , that is,*

$$0 \leq \xi_i \leq 1, \quad \xi_i \in C^\infty(V_i), \quad \text{supp } \xi_i \subset V_i, \quad \sum_{i=0}^k \xi_i = 1 \text{ in } \Omega. \quad (2.9)$$

Define

$$[f]^\varepsilon(x, t) := \xi_0(x)f^\varepsilon(x, t) + \sum_{i=1}^k \xi_i(x)\tilde{f}_i^\varepsilon(x, t), \quad \forall x \in \Omega, \quad (2.10)$$

where $f^\varepsilon(x, t)$ and $\tilde{f}_i^\varepsilon(x, t)$ are defined in (2.1) and (2.5) respectively.

Then, sending $\varepsilon \rightarrow 0$, we have

$$[f]^\varepsilon \rightarrow f \quad \text{in } L_{loc}^p(0, T; W^{1,p}(\Omega)). \quad (2.11)$$

Moreover,

$$[f]^\varepsilon - f^\varepsilon \rightarrow 0 \quad \text{in } L_{loc}^p(0, T; W^{1,p}(V_0)) \quad \text{and} \quad [f]^\varepsilon - \tilde{f}_i^\varepsilon \rightarrow 0 \quad \text{in } L_{loc}^p(0, T; W^{1,p}(V_i)). \quad (2.12)$$

Proof. . It follows from (2.7), (2.8), (2.9) that for any multi-index α satisfying $|\alpha| \leq 1$,

$$\begin{aligned} & \|\partial_x^\alpha([f]^\varepsilon - f)\|_{L_{loc}^p(0,T;L^p(\Omega))} \\ & \leq \|\partial_x^\alpha(\xi_0 f^\varepsilon - \xi_0 f)\|_{L_{loc}^p(0,T;L^p(V_0))} + \sum_{i=1}^k \|\partial_x^\alpha(\xi_i \tilde{f}_i^\varepsilon - \xi_i f)\|_{L_{loc}^p(0,T;L^p(V_i))} \\ & \leq C\|f^\varepsilon - f\|_{L_{loc}^p(0,T;W^{1,p}(V_0))} + C\sum_{i=1}^k \|\tilde{f}_i^\varepsilon - f\|_{L_{loc}^p(0,T;W^{1,p}(V_i))} \leq C\delta. \end{aligned}$$

This proves (2.11).

As a direct result of (2.11), (2.1) and (2.6), we obtain (2.12). \square

2.2. Commutator estimates.

Lemma 2.1 (Lemma 2.3 in [27]). *Suppose $\rho \in W^{1,r_1}([0, T] \times \Omega)$, $u \in L^{r_2}([0, T] \times \Omega)$, and $1 \leq r, r_1, r_2 \leq \infty$, $\frac{1}{r_1} + \frac{1}{r_2} = \frac{1}{r}$. Then,*

$$\|\partial(\rho u)^\varepsilon - \partial(\rho u^\varepsilon)\|_{L_{loc}^r([0,T] \times \Omega)} \leq \|u\|_{L^{r_2}([0,T] \times \Omega)} \|\partial\rho\|_{L^{r_1}([0,T] \times \Omega)}, \quad (2.13)$$

where $\partial = \partial_t$ or $\partial = \partial_x$, and f^ε is defined as in (2.1). Furthermore,

$$\partial(\rho u)^\varepsilon - \partial(\rho u^\varepsilon) \rightarrow 0 \text{ in } L_{loc}^r([0, T] \times \Omega), \quad \text{as } \varepsilon \rightarrow 0, \quad (2.14)$$

where $\underline{r} = r$ if $r_2 < \infty$ and $\underline{r} < r$ if $r_2 = \infty$.

We will need the following variant of Lemma 2.1.

Corollary 2.1. *Under the same assumptions listed in Lemma 2.1, we have*

$$\|\partial_t((\tilde{\rho}\tilde{u})_i^\varepsilon - \rho\tilde{u}_i^\varepsilon)\|_{L_{loc}^r(0, T; L^r(V_i))} \leq C\|u\|_{L^2([0, T] \times \Omega)} (\|\partial_t \rho\|_{L^1([0, T] \times \Omega)} + \|\nabla \rho\|_{L^1([0, T] \times \Omega)}) \quad (2.15)$$

and

$$\|\partial_x((\tilde{\rho}\tilde{u})_i^\varepsilon - \rho\tilde{u}_i^\varepsilon)\|_{L_{loc}^r(0, T; L^r(V_i))} \leq C\|u\|_{L^2([0, T] \times \Omega)} \|\nabla \rho\|_{L^1([0, T] \times \Omega)}, \quad (2.16)$$

where V_i ($i = 1, \dots, k$) is the same as mentioned earlier, and \tilde{f}_i^ε is defined in (2.5). Furthermore,

$$\partial((\tilde{\rho}\tilde{u})_i^\varepsilon - \rho\tilde{u}_i^\varepsilon) \rightarrow 0 \text{ in } L_{loc}^r(0, T; L^r(V_i)) \quad \text{as } \varepsilon \rightarrow 0, \quad (2.17)$$

where \underline{r} is given as in Lemma 2.1.

Proof. The proof is the same spirit of Lemma 2.1. By (2.13),

$$\begin{aligned} & \|\partial_t((\tilde{\rho}\tilde{u})_i^\varepsilon - \rho\tilde{u}_i^\varepsilon)\|_{L_{loc}^r(0, T; L^r(V_i))} \\ & \leq \|\partial_t((\tilde{\rho}\tilde{u})_i^\varepsilon - \tilde{\rho}\tilde{u}_i^\varepsilon)\|_{L_{loc}^r(0, T; L^r(V_i))} + \|\partial_t(\tilde{\rho} - \rho)\tilde{u}_i^\varepsilon\|_{L_{loc}^r(0, T; L^r(V_i))} \\ & \quad + \|(\tilde{\rho} - \rho)\partial_t\tilde{u}_i^\varepsilon\|_{L_{loc}^r(0, T; L^r(V_i))} \\ & \leq C\|u\|_{L^2([0, T] \times \Omega)} \|\partial_t \rho\|_{L^1([0, T] \times \Omega)} + \|(\tilde{\rho} - \rho)\partial_t\tilde{u}_i^\varepsilon\|_{L_{loc}^r(0, T; L^r(V_i))}. \end{aligned} \quad (2.18)$$

Thus, to prove (2.15), it suffices to estimate the last term in (2.18). Since

$$\begin{aligned} (\tilde{\rho} - \rho)\partial_t\tilde{u}_i^\varepsilon &= (\rho(x^\varepsilon, t) - \rho(x, t)) \int \int u(y, s) \partial_t \eta_\varepsilon(x^\varepsilon - y, t - s) dy ds \\ &\leq C \frac{|\rho(x^\varepsilon, t) - \rho(x, t)|}{\varepsilon} \int_{(t-\varepsilon, t+\varepsilon)} \int_{B(x^\varepsilon, \varepsilon)} \frac{|u(y, s)|}{\varepsilon^4} dy ds \\ &\leq C \left(\int_0^1 |\nabla \rho(x + \tau \varepsilon e_3, t)| d\tau \right) (|u| * \tilde{J}_\varepsilon) \end{aligned}$$

with $\tilde{J}_\varepsilon(x, t) = \frac{1}{\varepsilon^4} \mathbf{1}_{B(0, \varepsilon)}(x^\varepsilon, t)$, then,

$$\begin{aligned} \|(\tilde{\rho} - \rho)\partial_t\tilde{u}_i^\varepsilon\|_{L_{loc}^r(0, T; L^r(V_i))} &\leq C\|u\| * J_\varepsilon \|_{L_{loc}^{r_2}(0, T; L^{r_2}(V_i))} \|\nabla \rho\|_{L^1(V_i \times (0, T))} \\ &\leq C\|u\|_{L^2([0, T] \times \Omega)} \|\nabla \rho\|_{L^1([0, T] \times \Omega)}. \end{aligned} \quad (2.19)$$

Therefore, combining (2.18) with (2.19), we get (2.15).

The argument for (2.16) goes similarly, and (2.17) follows from (2.15), (2.16) by a density arguments. \square

Lemma 2.2. *Let $1 \leq r, r_1, r_2 < \infty$, $\frac{1}{r_1} + \frac{1}{r_2} = \frac{1}{r}$, and $f \in L^{r_1}$, $g \in L^{r_2}$. Then*

$$(fg)^\varepsilon - fg^\varepsilon \rightarrow 0 \text{ in } L_{loc}^r([0, T] \times \Omega) \quad \text{as } \varepsilon \rightarrow 0. \quad (2.20)$$

Proof. The Hölder inequality gives

$$\begin{aligned} |(fg)^\varepsilon - fg^\varepsilon| &= \left| \int \int (f(y, s) - f(x, t))g(y, s)\eta_\varepsilon(x - y, t - s)dyds \right| \\ &\leq \left(\frac{1}{\varepsilon^4} \int_{t-\varepsilon}^{t+\varepsilon} \int_{B(x, \varepsilon)} |f(x, t) - f(y, s)|^{r_1} \right)^{\frac{1}{r_1}} \left(\frac{1}{\varepsilon^4} \int_{t-\varepsilon}^{t+\varepsilon} \int_{B(x, \varepsilon)} |g|^{r_2} \right)^{\frac{1}{r_2}} \\ &\leq \left(\frac{1}{\varepsilon^4} \int_{B(x, \varepsilon)} |f(x, t) - f(y, s)|^{r_1} \right)^{\frac{1}{r_1}} (|g|^{r_2} * J_\varepsilon)^{\frac{1}{r_2}}, \end{aligned}$$

with $J_\varepsilon(x, t) = \frac{1}{\varepsilon^4} \mathbf{1}_{B(0, \varepsilon)}(x, t)$. The Lebesgue Differentiation Theorem implies that

$$\frac{1}{\varepsilon^4} \int_{t-\varepsilon}^{t+\varepsilon} \int_{B(x, \varepsilon)} |f(x, t) - f(y, s)|^{r_1} \rightarrow 0, \quad \text{as } \varepsilon \rightarrow 0.$$

Noticing that

$$\|(fg)^\varepsilon - fg^\varepsilon\|_{L^r_{loc}(\Omega \times (0, T))} \leq \|f\|_{L^{r_1}} \|g\|_{L^{r_2}} \leq C,$$

we obtain (2.20) by the Dominated Convergence Theorem. \square

The following Hardy-type imbedding will be useful for later use.

Lemma 2.3 ([21]). *Let $p \in [1, \infty)$ and $f \in W_0^{1,p}(\Omega)$. There is a constant C which depends on p and Ω , such that*

$$\left\| \frac{f(x)}{\text{dist}(x, \partial\Omega)} \right\|_{L^p(\Omega)} \leq C \|f\|_{W_0^{1,p}(\Omega)},$$

3. PROOF OF THEOREM 1.1

In the following, we will still use the conventions mentioned in Section 2.

Taking the j -th component of equations (1.1)₂, testing it against $\eta_\varepsilon(x^\varepsilon - y, t - s)$, summing up the expressions, and using (2.5), we deduce for every $(x, t) \in V_i \times (\varepsilon, T - \varepsilon)$ with $i \in \{1, 2, \dots, k\}$,

$$\partial_t(\tilde{\rho}\tilde{u})_i^\varepsilon + \text{div}(\tilde{\rho}\tilde{u} \otimes \tilde{u})_i^\varepsilon + \nabla(\tilde{\rho}^\gamma)_i^\varepsilon - \mu\Delta\tilde{u}_i^\varepsilon - (\mu + \lambda)\nabla\text{div}\tilde{u}_i^\varepsilon = 0. \quad (3.1)$$

To explain how we obtain (3.1), let us take the term $\text{div}(\tilde{\rho}\tilde{u} \otimes \tilde{u})_i^\varepsilon$ for example. In fact, by (2.5),

$$\begin{aligned} \text{div}(\tilde{\rho}\tilde{u} \otimes \tilde{u}^j)_i^\varepsilon(x, t) &= \int_0^T \int_{\mathcal{V}_i} \text{div}_y(\tilde{\rho}\tilde{u} \otimes \tilde{u}^j)(y, s)\eta_\varepsilon(x - y, t - s)dyds \\ &= \int_0^T \int_{\mathcal{V}_i} (\rho u \otimes u^j)(y - \varepsilon\tilde{n}(x_1), s) \cdot \nabla_x \eta_\varepsilon(x - y, t - s)dyds \\ &= \int_0^T \int_{\mathcal{V}_{i-\varepsilon\tilde{n}(x_1)}} \rho u \otimes u^j(y, s) \cdot \nabla_x \eta_\varepsilon(x^\varepsilon - y, t - s)dyds. \end{aligned}$$

Similarly, if we test (1.1)₂ against $\eta_\varepsilon(x - y, t - s)$, we infer that for every $(x, t) \in V_0 \times (\varepsilon, T - \varepsilon)$,

$$\partial_t(\rho u)^\varepsilon + \operatorname{div}(\rho u \otimes u)^\varepsilon + \nabla(\rho^\gamma)^\varepsilon - \mu \Delta u^\varepsilon - (\mu + \lambda) \nabla \operatorname{div} u^\varepsilon = 0. \quad (3.2)$$

Combining (3.1) with (3.2) implies that

$$\begin{aligned} & \partial_t \left(\xi_0(\rho u)^\varepsilon + \sum_{i=1}^k \xi_i(\tilde{\rho} \tilde{u})_i^\varepsilon \right) + \left(\xi_0 \operatorname{div}(\rho u \otimes u)^\varepsilon + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{\rho} \tilde{u} \otimes \tilde{u})_i^\varepsilon \right) \\ & + \left(\xi_0 \nabla(\rho^\gamma)^\varepsilon + \sum_{i=1}^k \xi_i \nabla(\tilde{\rho}^\gamma)_i^\varepsilon \right) \\ & - \left(\xi_0(\mu \Delta u^\varepsilon + (\mu + \lambda) \nabla \operatorname{div} u^\varepsilon) + \sum_{i=1}^k \xi_i(\mu \Delta \tilde{u}_i^\varepsilon + (\mu + \lambda) \nabla \operatorname{div} \tilde{u}_i^\varepsilon) \right) = 0, \end{aligned} \quad (3.3)$$

where ξ_0 and ξ_i are given in (2.9).

Next, we fix small constants $\tau > 0$, $\delta > 0$, and define the cut-off functions $\psi_\tau(t) \in C_0^1((\tau, T - \tau))$ and $\phi_\delta(x) \in C_0^1(\Omega)$ satisfying

$$\begin{cases} 0 \leq \phi_\delta(x) \leq 1, & \phi_\delta(x) = 1 \text{ if } x \in \Omega \text{ and } \operatorname{dist}(x, \partial\Omega) \geq \delta, \\ \phi_\delta \rightarrow 1 \text{ as } \delta \rightarrow 0, & \text{and } |\nabla \phi_\delta| \leq \frac{2}{\operatorname{dist}(x, \partial\Omega)}. \end{cases} \quad (3.4)$$

This way $\psi_\tau \phi_\delta [u]^\varepsilon$ is a legitimate test function, where $[u]^\varepsilon$ is defined in (2.10). Multiplying (3.3) by $\psi_\tau \phi_\delta [u]^\varepsilon$ and integrating it over $\Omega \times (0, T)$ leads to

$$\begin{aligned} & \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \partial_t \left(\xi_0(\rho u)^\varepsilon + \sum_{i=1}^k \xi_i(\tilde{\rho} \tilde{u})_i^\varepsilon \right) \\ & + \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \left(\xi_0 \operatorname{div}(\rho u \otimes u)^\varepsilon + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{\rho} \tilde{u} \otimes \tilde{u})_i^\varepsilon \right) \\ & + \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \left(\xi_0 \nabla(\rho^\gamma)^\varepsilon + \sum_{i=1}^k \xi_i \nabla(\tilde{\rho}^\gamma)_i^\varepsilon \right) \\ & - \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \left(\xi_0(\mu \Delta u^\varepsilon + (\mu + \lambda) \nabla \operatorname{div} u^\varepsilon) + \sum_{i=1}^k \xi_i(\mu \Delta \tilde{u}_i^\varepsilon + (\mu + \lambda) \nabla \operatorname{div} \tilde{u}_i^\varepsilon) \right) = 0. \end{aligned} \quad (3.5)$$

In the rest, we will calculate the terms in (3.5) one by one, and send ε , δ , τ to zero in the following three steps.

3.1. Step 1: ε -limit for (3.5).

Lemma 3.1. *For fixed τ and δ , the first two terms in (3.5) satisfy*

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t \left(\xi_0 (\rho u)^{\varepsilon} + \sum_{i=1}^k \xi_i (\tilde{\rho} \tilde{u})_i^{\varepsilon} \right) \\ & + \lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \operatorname{div}(\rho u \otimes u)^{\varepsilon} + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{\rho} \tilde{u} \otimes \tilde{u})_i^{\varepsilon} \right) \\ & = -\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \phi_{\delta} \rho |u|^2 - \frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |u|^2. \end{aligned} \quad (3.6)$$

Proof. Firstly, we have

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t \left(\xi_0 (\rho u)^{\varepsilon} + \sum_{i=1}^k \xi_i (\tilde{\rho} \tilde{u})_i^{\varepsilon} \right) \\ & = \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t \left(\xi_0 (\rho u)^{\varepsilon} + \sum_{i=1}^k \xi_i (\tilde{\rho} \tilde{u})_i^{\varepsilon} - \rho [u]^{\varepsilon} \right) + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t (\rho [u]^{\varepsilon}) \\ & =: I_1 + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t (\rho [u]^{\varepsilon}). \end{aligned} \quad (3.7)$$

Let us show

$$\lim_{\varepsilon \rightarrow 0} I_1 = 0. \quad (3.8)$$

The definition of $[u]^{\varepsilon}$ in (2.10) implies

$$\rho [u]^{\varepsilon} = \rho \left(\xi_0 u^{\varepsilon} + \sum_{i=1}^k \xi_i \tilde{u}_i^{\varepsilon} \right) = \xi_0 \rho u^{\varepsilon} + \sum_{i=1}^k \xi_i \rho \tilde{u}_i^{\varepsilon}.$$

This, along with (2.9), (2.15), (2.13), implies that

$$\begin{aligned} |I_1| & \leq C \int_{\tau}^{T-\tau} \|u\|_{L^q} \left(\|\xi_0 \partial_t((\rho u)^{\varepsilon} - \rho u^{\varepsilon})\|_{L^{\frac{q}{q-1}}(\Omega)} + \sum_{i=1}^k \|\xi_i \partial_t((\tilde{\rho} \tilde{u})_i^{\varepsilon} - \rho \tilde{u}_i^{\varepsilon})\|_{L^{\frac{q}{q-1}}(\Omega)} \right) \\ & \leq C \int_{\tau}^{T-\tau} \|u\|_{L^q} \left(\|\partial_t((\rho u)^{\varepsilon} - \rho u^{\varepsilon})\|_{L^{\frac{q}{q-1}}(V_0)} + \sum_{i=1}^k \|\partial_t((\tilde{\rho} \tilde{u})_i^{\varepsilon} - \rho \tilde{u}_i^{\varepsilon})\|_{L^{\frac{q}{q-1}}(V_i)} \right) \\ & \leq C \int_0^T \|u\|_{L^q}^2 \left(\|\partial_t \rho\|_{L^{\frac{q}{q-2}}} + \|\nabla \rho\|_{L^{\frac{q}{q-2}}} \right). \end{aligned} \quad (3.9)$$

On the other hand, it follows from (1.5), (1.7), (1.8) that

$$\rho_t = -(\rho \operatorname{div} u + 2\sqrt{\rho} u \cdot \nabla \sqrt{\rho}) \in L^2(0, T; L^2) + L^p(0, T; L^{\frac{2q}{q+2}}). \quad (3.10)$$

Therefore, from (3.10), (1.7), (1.8) we deduce

$$|I_1| \leq C \int_0^T \left(\|u\|_{L^q}^p + \|\partial_t \rho\|_{L^{\frac{q}{q-2}}}^{\frac{p}{p-2}} + \|\nabla \rho\|_{L^{\frac{q}{q-2}}}^{\frac{p}{p-2}} \right) \leq C,$$

provided that $p \geq 4$, $q \geq 6$.

Furthermore, with (3.10) and (1.7), from Lemma 2.1 and Corollary 2.1 we obtain

$$\begin{aligned} \partial_t ((\rho u)^\varepsilon - \rho u^\varepsilon) &\text{ in } L_{loc}^{\frac{2p}{p+2}} \left(0, T; L^{\frac{2q}{q+4}}(V_0) \right), \\ \partial_t ((\tilde{\rho}\tilde{u})_i^\varepsilon - \rho\tilde{u}_i^\varepsilon) &\rightarrow 0 \text{ in } L_{loc}^{\frac{2p}{p+2}} \left(0, T; L^{\frac{2q}{q+4}}(V_i) \right). \end{aligned}$$

This concludes (3.8).

Secondly, the convection term can be treated as

$$\begin{aligned} &\int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \left(\xi_0 \operatorname{div}(\rho u \otimes u)^\varepsilon + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{\rho}\tilde{u} \otimes \tilde{u})_i^\varepsilon \right) \\ &= \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \left(\xi_0 \operatorname{div}(\rho u \otimes u)^\varepsilon + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{\rho}\tilde{u} \otimes \tilde{u})_i^\varepsilon - \operatorname{div}(\rho u \otimes [u]^\varepsilon) \right) \\ &\quad + \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \operatorname{div}(\rho u \otimes [u]^\varepsilon) \\ &=: I_2 + \int_0^T \int_\Omega \psi_\tau \phi_\delta [u]^\varepsilon \operatorname{div}(\rho u \otimes [u]^\varepsilon). \end{aligned} \tag{3.11}$$

We claim that

$$\lim_{\varepsilon \rightarrow 0} I_2 = 0. \tag{3.12}$$

In fact, by (2.9),

$$\begin{aligned} |I_2| &\leq C \int_\tau^{T-\tau} \|u\|_{L^q} \|\operatorname{div}((\rho u \otimes u)^\varepsilon - \rho u \otimes [u]^\varepsilon)\|_{L^{\frac{q}{q-1}}(V_0)} \\ &\quad + C \sum_{i=1}^k \int_\tau^{T-\tau} \|u\|_{L^q} \|\operatorname{div}((\tilde{\rho}\tilde{u} \otimes \tilde{u})_i^\varepsilon - \rho u \otimes [u]^\varepsilon)\|_{L^{\frac{q}{q-1}}(V_i)} \\ &=: I_{21} + \sum_{i=1}^k I_{22}^i. \end{aligned} \tag{3.13}$$

Making use of (2.13) and (1.7), one has

$$\|\nabla(\rho u)\|_{L^{\frac{q}{q-2}}} \leq C \|u\|_{H^1} (1 + \|\nabla\sqrt{\rho}\|_{L^2}) \leq C \|\nabla u\|_{L^2}.$$

Thus,

$$\begin{aligned} &|I_{21}| \\ &\leq \int_\tau^{T-\tau} \|u\|_{L^q} \left(\|\operatorname{div}((\rho u \otimes u)^\varepsilon - \rho u \otimes u^\varepsilon)\|_{L^{\frac{q}{q-1}}(V_0)} + \|\operatorname{div}(\rho u \otimes (u^\varepsilon - [u]^\varepsilon))\|_{L^{\frac{q}{q-1}}(V_0)} \right) \\ &\leq C \int_0^T \|u\|_{L^q}^2 \|\nabla(\rho u)\|_{L^{\frac{q}{q-2}}} + \int_\tau^{T-\tau} \|u\|_{L^q} \|\operatorname{div}(\rho u \otimes (u^\varepsilon - [u]^\varepsilon))\|_{L^{\frac{q}{q-1}}(V_0)} \\ &\leq C \int_0^T (\|u\|_{L^q}^4 + \|\nabla u\|_{L^2}^2) + \int_\tau^{T-\tau} \|u\|_{L^q} \|\operatorname{div}(\rho u \otimes (u^\varepsilon - [u]^\varepsilon))\|_{L^{\frac{q}{q-1}}(V_0)}, \end{aligned} \tag{3.14}$$

Notice that

$$\begin{aligned}
& \int_{\tau}^{T-\tau} \|u\|_{L^q} \|\operatorname{div}(\rho u \otimes (u^\varepsilon - [u]^\varepsilon))\|_{L^{\frac{q}{q-1}}(V_0)} \\
& \leq C \int_{\tau}^{T-\tau} \|u\|_{L^q} \left(\|\nabla(\rho u)\|_{L^{\frac{q}{q-2}}} \|u^\varepsilon - [u]^\varepsilon\|_{L^q(V_0)} + \|\rho u\|_{L^q} \|\nabla(u^\varepsilon - [u]^\varepsilon)\|_{L^{\frac{q}{q-2}}(V_0)} \right) \\
& \leq C \int_{\tau}^{T-\tau} \|u\|_{L^q} (\|u\|_{H^1} \|u^\varepsilon - [u]^\varepsilon\|_{L^q(V_0)} + \|\nabla(u^\varepsilon - [u]^\varepsilon)\|_{L^2(V_0)} \|u\|_{L^q}) \\
& \leq C \left(\int_{\tau}^{T-\tau} \|u^\varepsilon - [u]^\varepsilon\|_{L^q(V_0)}^p + \|\nabla(u^\varepsilon - [u]^\varepsilon)\|_{L^2(V_0)}^2 \right)^{\frac{1}{2}},
\end{aligned} \tag{3.15}$$

and, owing to (1.5), (1.8), (2.12),

$$\lim_{\varepsilon \rightarrow 0} \left(\|u^\varepsilon - [u]^\varepsilon\|_{L^p_{loc}(0,T;L^q(V_0))} + \|\nabla(u^\varepsilon - [u]^\varepsilon)\|_{L^2_{loc}(0,T;L^2(V_0))} \right) = 0. \tag{3.16}$$

We conclude from (3.16) and (2.14) that

$$\lim_{\varepsilon \rightarrow 0} I_{21} = 0. \tag{3.17}$$

By (2.16), a similar argument to (3.14) and (3.15) infers that

$$\begin{aligned}
& |I_{22}^i| \\
& \leq \int_{\tau}^{T-\tau} \|u\|_{L^q} \left(\|\operatorname{div}((\tilde{\rho} \tilde{u} \otimes \tilde{u})_i^\varepsilon - \rho u \otimes \tilde{u}_i^\varepsilon)\|_{L^{\frac{q}{q-1}}(V_i)} + \|\operatorname{div}(\rho u \otimes (\tilde{u}_i^\varepsilon - [u]^\varepsilon))\|_{L^{\frac{q}{q-1}}(V_i)} \right) \\
& \leq C \int_0^T (\|u\|_{L^q}^4 + \|\nabla u\|_{L^2}^2) + C \int_{\tau}^{T-\tau} \left(\|u^\varepsilon - [u]^\varepsilon\|_{L^q(V_i)}^p + \|\nabla(u^\varepsilon - [u]^\varepsilon)\|_{L^2(V_i)}^2 \right).
\end{aligned} \tag{3.18}$$

and consequently, from (2.17) and (3.16) we get

$$\lim_{\varepsilon \rightarrow 0} I_{22}^i = 0.$$

This together with (3.17) implies (3.12).

Next, by the continuity equation (1.1)₁, a simple computation shows that

$$\begin{aligned}
& \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^\varepsilon \partial_t(\rho [u]^\varepsilon) + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^\varepsilon \operatorname{div}(\rho u \otimes [u]^\varepsilon) \\
& = -\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \phi_{\delta} \rho |[u]^\varepsilon|^2 - \frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |[u]^\varepsilon|^2 \\
& \quad - \int_0^T \int_{\Omega} \rho \partial_t \left(\phi_{\delta} \psi_{\tau} \frac{|[u]^\varepsilon|^2}{2} \right) - \rho u \cdot \nabla \left(\phi_{\delta} \psi_{\tau} \frac{|[u]^\varepsilon|^2}{2} \right) \\
& = -\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \phi_{\delta} \rho |[u]^\varepsilon|^2 - \frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |[u]^\varepsilon|^2 \\
& \rightarrow -\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \phi_{\delta} \rho |u|^2 - \frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |u|^2, \quad \text{as } \varepsilon \rightarrow 0.
\end{aligned} \tag{3.19}$$

In conclusion, owing to (3.8), (3.12), (3.19), we get (3.6) from (3.7) and (3.11). \square

Lemma 3.2. *For fixed τ and δ , the pressure term in (3.5) satisfies*

$$\lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla (\rho^{\gamma})^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla (\tilde{\rho}^{\gamma})_i^{\varepsilon} \right) = \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} u \cdot \nabla \rho^{\gamma}. \quad (3.20)$$

Proof. Owing to (1.7), we have

$$\nabla \rho^{\gamma} \in L^2(0, T; L^2). \quad (3.21)$$

We write

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla (\rho^{\gamma})^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla (\tilde{\rho}^{\gamma})_i^{\varepsilon} \right) \\ &= \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla (\rho^{\gamma})^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla (\tilde{\rho}^{\gamma})_i^{\varepsilon} - \nabla \rho^{\gamma} \right) + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \cdot \nabla \rho^{\gamma} \\ &=: I_3 + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \cdot \nabla \rho^{\gamma}, \end{aligned} \quad (3.22)$$

where the last integral makes sense due to (1.5) and (3.21).

Observe from (2.1) and (2.6) that

$$\lim_{\varepsilon \rightarrow 0} \left(\|\nabla((\rho^{\gamma})^{\varepsilon} - \rho^{\gamma})\|_{L^2_{loc}(0, T; L^2(V_0))} + \|\nabla((\tilde{\rho}^{\gamma})_i^{\varepsilon} - \rho^{\gamma})\|_{L^2_{loc}(0, T; L^2(V_i))} \right) = 0, \quad (3.23)$$

and hence,

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} |I_3| \\ & \leq C \lim_{\varepsilon \rightarrow 0} \|u\|_{L^2(0, T; L^2(\Omega))} \left(\|\xi_0 \nabla((\rho^{\gamma})^{\varepsilon} - \rho^{\gamma}) + \sum_{i=1}^k \xi_i \nabla((\tilde{\rho}^{\gamma})_i^{\varepsilon} - \rho^{\gamma})\|_{L^2_{loc}(0, T; L^2(\Omega))} \right) \\ & \leq C \lim_{\varepsilon \rightarrow 0} \left(\|\nabla((\rho^{\gamma})^{\varepsilon} - \rho^{\gamma})\|_{L^2_{loc}(0, T; L^2(V_0))} + \|\nabla((\tilde{\rho}^{\gamma})_i^{\varepsilon} - \rho^{\gamma})\|_{L^2_{loc}(0, T; L^2(V_i))} \right) \\ & = 0. \end{aligned} \quad (3.24)$$

Taking (3.24), (2.11) into accounts, we take $\varepsilon \rightarrow 0$ in (3.22) and complete the proof for Lemma 3.2. \square

Lemma 3.3. *For fixed τ and δ , the diffusion terms in (3.5) satisfy*

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 (\mu \Delta u^{\varepsilon} + (\mu + \lambda) \nabla \operatorname{div} u^{\varepsilon}) + \sum_{i=1}^k \xi_i (\mu \Delta \tilde{u}_i^{\varepsilon} + (\mu + \lambda) \nabla \operatorname{div} \tilde{u}_i^{\varepsilon}) \right) \\ &= - \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} (\mu |\nabla u|^2 + (\mu + \lambda) (\operatorname{div} u)^2) - \int_0^T \int_{\Omega} \psi_{\tau} \nabla \phi_{\delta} (\mu u \nabla u + (\mu + \lambda) u \operatorname{div} u). \end{aligned} \quad (3.25)$$

Proof. We see that

$$\begin{aligned}
& \mu \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \Delta u^{\varepsilon} + \sum_{i=1}^k \xi_i \Delta \tilde{u}^{\varepsilon} \right) \\
&= \mu \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \Delta u^{\varepsilon} + \sum_{i=1}^k \xi_i \Delta \tilde{u}^{\varepsilon} - \Delta [u]^{\varepsilon} \right) + \mu \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \Delta [u]^{\varepsilon} \quad (3.26) \\
&=: I_4 - \mu \int_0^T \int_{\Omega} \psi_{\tau} \nabla \phi_{\delta} [u]^{\varepsilon} \nabla [u]^{\varepsilon} - \mu \int_0^T \int_{\Omega} \phi_{\delta} \psi_{\tau} \nabla [u]^{\varepsilon} : \nabla [u]^{\varepsilon}.
\end{aligned}$$

We compute

$$\begin{aligned}
I_4 &= \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \Delta (u^{\varepsilon} - [u]^{\varepsilon}) + \sum_{i=1}^k \xi_i \Delta (\tilde{u}_i^{\varepsilon} - [u]^{\varepsilon}) \right) \\
&= - \int_0^T \int_{\Omega} \psi_{\tau} (\nabla \phi_{\delta} [u]^{\varepsilon} \xi_0 + \phi_{\delta} \operatorname{div} [u]^{\varepsilon} \xi_0 + \phi_{\delta} [u]^{\varepsilon} \nabla \xi_0) \nabla (u^{\varepsilon} - [u]^{\varepsilon}) \\
&\quad - \sum_{i=1}^k \int_0^T \int_{\Omega} \psi_{\tau} (\nabla \phi_{\delta} [u]^{\varepsilon} \xi_i + \phi_{\delta} \operatorname{div} [u]^{\varepsilon} \xi_i + \phi_{\delta} [u]^{\varepsilon} \nabla \xi_i) \nabla (\tilde{u}_i^{\varepsilon} - [u]^{\varepsilon}).
\end{aligned}$$

Thus

$$|I_4| \leq C(\delta) \left(\|\nabla u^{\varepsilon} - \nabla [u]^{\varepsilon}\|_{L^2_{loc}(0,T;L^2(V_0))} + \sum_{i=1}^k \|\nabla \tilde{u}_i^{\varepsilon} - \nabla [u]^{\varepsilon}\|_{L^2_{loc}(0,T;L^2(V_i))} \right).$$

In view of (1.5), (2.11), (2.12), it yields from (3.26) that

$$\begin{aligned}
& \mu \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \Delta u^{\varepsilon} + \sum_{i=1}^k \xi_i \Delta u^{\varepsilon} \right) \\
&\rightarrow -\mu \int_0^T \int_{\Omega} \psi_{\tau} \nabla \phi_{\delta} u \nabla u - \mu \int_0^T \int_{\Omega} \phi_{\delta} \psi_{\tau} \nabla u : \nabla u,
\end{aligned}$$

as $\varepsilon \rightarrow 0$. Applying the similar arguments for other diffusion terms, we get (3.25), and hence the lemma is proved. \square

In summary, Lemma 3.1-3.3 and equality (3.5) imply that

$$\begin{aligned}
& -\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \phi_{\delta} \rho |u|^2 - \frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |u|^2 \\
& - \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} \rho^{\gamma} \operatorname{div} u - \int_0^T \int_{\Omega} \psi_{\tau} u \cdot \nabla \phi_{\delta} \rho^{\gamma} \\
& - \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} (\mu |\nabla u|^2 + (\mu + \lambda) (\operatorname{div} u)^2) - \int_0^T \int_{\Omega} \psi_{\tau} \nabla \phi_{\delta} (\mu u \nabla u + (\mu + \lambda) u \operatorname{div} u) = 0. \quad (3.27)
\end{aligned}$$

Next we consider taking the δ -limit.

3.2. Step 2: δ -limit for (3.27). Thanks to (1.5), (1.2), (1.7), (1.8), and Lemma 2.3, it follows that

$$\begin{aligned}
& -\frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} \rho u \cdot \nabla \phi_{\delta} |u|^2 - \int_0^T \int_{\Omega} \psi_{\tau} u \cdot \nabla \phi_{\delta} \rho^{\gamma} - \int_0^T \int_{\Omega} \psi_{\tau} \nabla \phi_{\delta} (\mu u \nabla u + (\mu + \lambda) u \operatorname{div} u) \\
& \leq C \| |u| |\nabla \phi_{\delta}| \|_{L^2(0,T;L^2)} \left(\int_0^T \int_{\{x: \operatorname{dist}(x, \partial\Omega) < \delta\}} |u|^4 + \rho^2 + |\nabla u|^2 \right)^{\frac{1}{2}} \\
& \leq C \left\| \frac{u}{\operatorname{dist}(x, \partial\Omega)} \right\|_{L^2(0,T;L^2)} \left(\int_0^T \int_{\{x: \operatorname{dist}(x, \partial\Omega) < \delta\}} |u|^4 + \rho^2 + |\nabla u|^2 \right)^{\frac{1}{2}} \\
& \rightarrow 0, \quad \text{as } \delta \rightarrow 0.
\end{aligned} \tag{3.28}$$

By (3.28), taking $\delta \rightarrow 0$ in (3.27),

$$\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \rho |u|^2 + \int_0^T \int_{\Omega} \psi_{\tau} \rho^{\gamma} \operatorname{div} u + \int_0^T \int_{\Omega} \psi \nabla u : \mathbb{S} = 0. \tag{3.29}$$

On the other hand, it follows from (1.1)₁, (1.5), (1.7) that

$$\int_0^T \int_{\Omega} \psi_{\tau} \rho^{\gamma} \operatorname{div} u = \int_0^T \int_{\Omega} \psi_{\tau} \rho^{\gamma-1} (\rho_t + u \cdot \nabla \rho) = \frac{1}{\gamma-1} \int_0^T \int_{\Omega} \psi'_{\tau} \rho^{\gamma}.$$

Thus, (3.29) becomes

$$\frac{1}{2} \int_0^T \int_{\Omega} \psi'_{\tau} \rho |u|^2 + \frac{1}{\gamma-1} \int_0^T \int_{\Omega} \psi'_{\tau} \rho^{\gamma} + \int_0^T \int_{\Omega} \psi_{\tau} (\mu |\nabla u|^2 + (\mu + \lambda) |\operatorname{div} u|^2) = 0. \tag{3.30}$$

Denote

$$E(t) := \int_{\Omega} \left(\frac{1}{2} \rho |u|^2 + \frac{\rho^{\gamma}}{\gamma-1} \right) dx, \quad D(t) := \int_0^t \int_{\Omega} (\mu |\nabla u|^2 + (\mu + \lambda) |\operatorname{div} u|^2) dx ds. \tag{3.31}$$

Then (3.30) implies that

$$(E - D)' = 0, \quad \text{in } \mathcal{D}'((0, T)).$$

3.3. Step 3: Global energy balance. Finally we will obtain the exact energy equality on the whole time interval $[0, T]$. First we note that $D \in C([0, T])$. Second, we see that an approximation argument shows that (3.30) remains valid for functions ψ_{τ} belonging only to $W^{1,\infty}$ rather than C^1 .

It follows from (1.1)₁ that for any $\alpha \geq \frac{1}{2}$,

$$\partial_t(\rho^{\alpha}) = -\alpha \rho^{\alpha} \operatorname{div} u - 2\alpha \rho^{\alpha-\frac{1}{2}} u \cdot \nabla \sqrt{\rho},$$

which, together with (1.5) and (1.7), implies

$$\rho^{\alpha} \in L^{\infty}(0, T; H^1(\Omega)), \quad \partial_t(\rho^{\alpha}) \in L^2(0, T; L^{\frac{3}{2}}(\Omega)).$$

Hence, by the Aubin–Lions Lemma (cf. [34]),

$$\rho^{\alpha} \in C([0, T]; L^r(\Omega)), \quad (r < 6).$$

In particular we know that

$$\rho^{\gamma} \in C([0, T]; L^1(\Omega)). \tag{3.32}$$

In a similar way,

$$\rho u \in L^\infty(0, T; L^2(\Omega)) \cap H^1(0, T; W^{-1,1}(\Omega)) \hookrightarrow C([0, T]; L^2_{weak}(\Omega)). \quad (3.33)$$

Recalling (1.6) and the (3.32), we have

$$\begin{aligned} 0 &\leq \overline{\lim}_{t \rightarrow 0} \int_{\Omega} |\sqrt{\rho}u - \sqrt{\rho_0}u_0|^2 dx \\ &= 2\overline{\lim}_{t \rightarrow 0} \left(\int_{\Omega} \left(\frac{1}{2}\rho|u|^2 + \frac{1}{\gamma-1}\rho^\gamma \right) dx - \int_{\Omega} \left(\frac{1}{2}\rho_0|u_0|^2 + \frac{1}{\gamma-1}\rho_0^\gamma \right) dx \right) \\ &\quad + 2\overline{\lim}_{t \rightarrow 0} \left(\int_{\Omega} \sqrt{\rho_0}u_0(\sqrt{\rho_0}u_0 - \sqrt{\rho}u) dx + \frac{1}{\gamma-1} \int_{\Omega} (\rho_0^\gamma - \rho^\gamma) dx \right) \\ &\leq 2\overline{\lim}_{t \rightarrow 0} \int_{\Omega} \sqrt{\rho_0}u_0(\sqrt{\rho_0}u_0 - \sqrt{\rho}u) dx, \end{aligned} \quad (3.34)$$

and furthermore,

$$\begin{aligned} \overline{\lim}_{t \rightarrow 0} \int_{\Omega} \sqrt{\rho_0}u_0(\sqrt{\rho_0}u_0 - \sqrt{\rho}u) dx &= \overline{\lim}_{t \rightarrow 0} \int_{\Omega} u_0(\rho_0 u_0 - \rho u) dx + \overline{\lim}_{t \rightarrow 0} \int_{\Omega} u_0 \sqrt{\rho}u(\sqrt{\rho} - \sqrt{\rho_0}) dx \\ &= 0, \end{aligned} \quad (3.35)$$

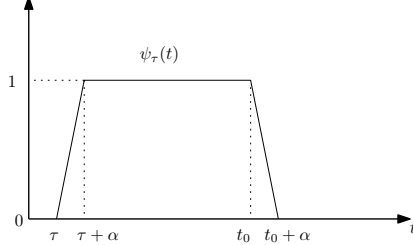
owing to (3.33), (3.32) and (1.9). Therefore, from (3.34), (3.35), and (3.33) we deduce that

$$(\sqrt{\rho}u)(t) \rightarrow (\sqrt{\rho}u)(0) \quad \text{strongly in } L^2(\Omega) \text{ as } t \rightarrow 0^+. \quad (3.36)$$

Similarly, one has the right temporal continuity of $\sqrt{\rho}u$ in L^2 , that is, for any $t_0 \geq 0$,

$$(\sqrt{\rho}u)(t) \rightarrow (\sqrt{\rho}u)(t_0) \quad \text{strongly in } L^2(\Omega) \text{ as } t \rightarrow t_0^+. \quad (3.37)$$

Now for $t_0 > 0$, we choose some positive τ and α such that $\tau + \alpha < t_0$ and define the following test function

$$\psi_\tau(t) = \begin{cases} 0, & 0 \leq t \leq \tau, \\ \frac{t-\tau}{\alpha}, & \tau \leq t \leq \tau + \alpha, \\ 1, & \tau + \alpha \leq t \leq t_0, \\ \frac{t_0-t}{\alpha}, & t_0 \leq t \leq t_0 + \alpha, \\ 0, & t_0 + \alpha \leq t. \end{cases}$$


Substituting the above test function into (3.30) we obtain that

$$\begin{aligned} &\frac{1}{\alpha} \int_{\tau}^{\tau+\alpha} \int_{\Omega} \left(\frac{1}{2}\rho|u|^2 + \frac{\rho^\gamma}{\gamma-1} \right) dx ds - \frac{1}{\alpha} \int_{t_0}^{t_0+\alpha} \int_{\Omega} \left(\frac{1}{2}\rho|u|^2 + \frac{\rho^\gamma}{\gamma-1} \right) dx ds \\ &= - \int_{\tau}^{t_0+\alpha} \int_{\Omega} \psi_\tau (\mu|\nabla u|^2 + (\mu + \lambda)|\operatorname{div}u|^2) dx ds = D(\tau) - D(t_0 + \alpha), \end{aligned}$$

where D is defined in (3.31).

Sending $\alpha \rightarrow 0$, using the right continuity of $\sqrt{\rho}u$ in L^2 and the continuity of ρ^γ in L^1 (cf. (3.32) and (3.37)), and the continuity of $D(t)$, one has

$$E(\tau) - E(t_0) = D(\tau) - D(t_0).$$

Finally sending $\tau \rightarrow 0$, from (3.36) we have

$$(E - D)(t_0) = (E - D)(0),$$

which is exactly (1.4), and hence we complete the proof of Theorem 1.1.

APPENDIX A. APPLICATION TO THE INCOMPRESSIBLE NAVIER–STOKES EQUATIONS

In this appendix, we apply our global approximation method to the incompressible Navier–Stokes equations posed on a bounded domain:

$$\begin{cases} u_t + \operatorname{div}(u \otimes u) + \nabla P - \nu \Delta u = 0, \\ \operatorname{div} u = 0, \\ u|_{\partial\Omega} = 0, \\ u(x, 0) = u_0(x), \quad \operatorname{div} u_0 = 0. \end{cases} \quad (\text{A.1})$$

It is well-known that the Leray–Hopf weak solution u to (A.1) satisfies

$$u \in L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H_0^1(\Omega)) \quad (\text{A.2})$$

and the following energy inequality

$$\int_\Omega \frac{1}{2}|u|^2 dx + \nu \int_0^t \int_\Omega |\nabla u|^2 dx ds \leq \int_\Omega \frac{1}{2}|u_0|^2 dx. \quad (\text{A.3})$$

We first recall a result of [29, Theorem 1] regarding the pressure field associated to the Leray–Hopf solution of (A.1).

Theorem A.1. (*[29, Theorem 1]*) *Assume that Ω is an open, bounded domain with C^2 boundary $\partial\Omega$, and u is a Leray–Hopf solution of (A.1). Then there exists a pressure field $P \in L^r(0, T; W^{1,s}(\Omega))$ with*

$$\frac{3}{s} + \frac{2}{r} = 4, \quad \frac{4}{3} < s < \frac{3}{2}, \quad (\text{A.4})$$

such that for all $\varphi \in C_0^\infty((0, T) \times \Omega)$,

$$\int_0^T \int_\Omega (u \cdot \partial_t \varphi + u \otimes u : \nabla \varphi + P \operatorname{div} \varphi + \nu u \cdot \Delta \varphi) dx dt = 0.$$

Remark A.1. The additional smoothness of the boundary is assumed to ensure the existence of the Leray–Hopf weak solution of (A.1).

Remark A.2. Another important implication of Theorem A.1 is that it allows us to use the test function $\psi_\tau \phi_\delta[u]^\varepsilon$ we introduced in the previous sections which is not solenoidal to test against the incompressible Navier–Stokes equations.

The main result concerning energy equality of (A.1) is

Theorem A.2. *Assume that Ω is an open, bounded domain with C^2 boundary $\partial\Omega$, and u is a Leray–Hopf weak solution of (A.1). Then the equality in (A.3) is achieved, provided that*

$$u \in L^p(0, T; L^q(\Omega)), \quad \frac{1}{p} + \frac{1}{q} \leq \frac{1}{2}, \quad 4 \leq q, \quad (\text{A.5})$$

and the associated pressure given in Theorem A.1 satisfies

$$P \in L^2(0, T; L^2(\partial\Omega)). \quad (\text{A.6})$$

Remark A.3. A notable difference between our Theorem A.1 and the result in [37], as indicated in the Introduction, is that we do not need to assume any additional Besov regularity on the velocity to handle the boundary effects coming from the diffusion term.

Remark A.4. From Theorem A.1 we know that the pressure field enjoys sufficient regularity to define its trace on the boundary $P \in W^{1-\frac{1}{s},s}(\partial\Omega) \subset L^s(\partial\Omega)$ for a.e. $t \in [0, T]$. In fact the fractional Sobolev embedding (see, e.g. [30]) further implies that

$$P \in L^{\tilde{s}}(\partial\Omega) \quad \text{where} \quad s \leq \tilde{s} \leq \frac{2s}{3-s}.$$

From (A.4) we see that $\tilde{s} < 2$. Here we need to assume a bit more integrability of the pressure trace (cf. (A.6)).

Remark A.5. Note that by interpolation we see that $L^2H^1 \cap L^4L^4$ lands in the Onsager-critical Besov spaces $L^3B_{3,r}^{1/3}$ for $1 \leq r < \infty$. It would be interesting to obtain the energy equality for velocities in the Onsager-critical Besov spaces $L^3B_{3,\infty}^{1/3}$ in the interior, a la Constantin et al. [6].

The proof of Theorem A.2 is a slight modification of that in Theorem 1.1. We only prove the Lemmas A.1–A.2 below to address conditions (A.5)–(A.6) and the main differences. An important ingredient in the argument is the global L^p estimate of the pressure, which is given in the following proposition.

Proposition A.1. *Let the assumptions of Theorem A.2 hold, then the pressure field satisfies*

$$\|P\|_{L^2(\Omega)} \leq C \left(\|u\|_{L^4(\Omega)}^2 + \|P|_{\partial\Omega}\|_{L^2} \right). \quad (\text{A.7})$$

Proof. The pressure satisfies a Poisson problem together with certain boundary regularity.

$$\begin{aligned} -\Delta P &= \operatorname{div}\operatorname{div}(u \otimes u) \quad \text{in } \Omega, \\ P|_{\partial\Omega} &\in L^2. \end{aligned}$$

Using duality and the method of transposition (e.g. [15, Lemma 2]) we see that

$$\|P\|_{L^2(\Omega)} \leq C \left(\|(u \otimes u)\|_{L^2(\Omega)} + \|P|_{\partial\Omega}\|_{L^2} \right) \leq C \left(\|u\|_{L^4(\Omega)}^2 + \|P|_{\partial\Omega}\|_{L^2} \right),$$

completing the proof of the proposition. \square

Let u be a Leray–Hopf weak solution to (A.1). The same deduction as (3.5) yields

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \partial_t \left(\xi_0 u^{\varepsilon} + \sum_{i=1}^k \xi_i \tilde{u}_i^{\varepsilon} \right) \\ & + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \operatorname{div}(u \otimes u)^{\varepsilon} + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{u} \otimes \tilde{u}_i)^{\varepsilon} \right) \\ & + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla P^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla \tilde{P}_i^{\varepsilon} \right) \\ & - \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \Delta u^{\varepsilon} + \sum_{i=1}^k \xi_i \Delta \tilde{u}_i^{\varepsilon} \right) = 0. \end{aligned} \quad (\text{A.8})$$

Lemma A.1. *The convection term in (A.8) satisfies*

$$\lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \operatorname{div}(u \otimes u)^{\varepsilon} + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{u} \otimes \tilde{u})_i^{\varepsilon} \right) = 0. \quad (\text{A.9})$$

Proof. A careful computation shows

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \operatorname{div}(u \otimes u)^{\varepsilon} + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{u} \otimes \tilde{u})_i^{\varepsilon} - \operatorname{div}(u \otimes [u]^{\varepsilon}) \right) \\ & \leq \int_{\tau}^{T-\tau} \|u\|_{L^4} \left(\|\operatorname{div}(u \otimes u)^{\varepsilon} - \operatorname{div}(u \otimes u^{\varepsilon})\|_{L^{\frac{4}{3}}(V_0)} + \|\operatorname{div}(u \otimes (u^{\varepsilon} - [u]^{\varepsilon}))\|_{L^{\frac{4}{3}}(V_0)} \right) \\ & \quad + \sum_{i=1}^k \int_{\tau}^{T-\tau} \|u\|_{L^4} \left(\|\operatorname{div}(\tilde{u} \otimes \tilde{u})_i^{\varepsilon} - \operatorname{div}(u \otimes \tilde{u}_i^{\varepsilon})\|_{L^{\frac{4}{3}}(V_i)} + \|\operatorname{div}(u \otimes (\tilde{u}_i^{\varepsilon} - [u]^{\varepsilon}))\|_{L^{\frac{4}{3}}(V_i)} \right) \\ & \leq C \int_0^T \|u\|_{L^4}^2 \|\nabla u\|_{L^2} \\ & \quad + C \int_{\tau}^{T-\tau} \|u\|_{L^4} (\|\nabla u\|_{L^2} \|u^{\varepsilon} - [u]^{\varepsilon}\|_{L^4(V_0)} + \|u\|_{L^4} \|\nabla(u^{\varepsilon} - [u]^{\varepsilon})\|_{L^4(V_0)}) \\ & \quad + C \sum_{i=1}^k \int_{\tau}^{T-\tau} \|u\|_{L^4} \left(\|\nabla u\|_{L^2} \|\tilde{u}_i^{\varepsilon} - [u]^{\varepsilon}\|_{L^{\frac{2q}{q-2}}(V_i)} + \|u\|_{L^4} \|\nabla(\tilde{u}_i^{\varepsilon} - [u]^{\varepsilon})\|_{L^4(V_i)} \right) \\ & \leq C \int_0^T (\|\nabla u\|_{L^2}^2 + \|u\|_{L^4}^4), \end{aligned}$$

owing to Lemma 2.1 and Corollary 2.1.

By virtue of (A.2) and (A.5), it satisfies

$$\|u\|_{L^4(0,T;L^4)} \leq \|u\|_{L^{\infty}(0,T;L^2)}^{\frac{q-4}{2q-4}} \|u\|_{L^{\frac{4q}{2q-4}}(0,T;L^q)}^{\frac{q}{2q-4}} \leq C, \quad (\text{A.10})$$

provided $q \geq 4$, $\frac{4q}{2q-4} \leq p$, which is equivalent to $q \geq 4$, $\frac{1}{q} + \frac{1}{p} \leq \frac{1}{2}$. Thus,

$$\lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \operatorname{div}(u \otimes u)^{\varepsilon} + \sum_{i=1}^k \xi_i \operatorname{div}(\tilde{u} \otimes \tilde{u})_i^{\varepsilon} - \operatorname{div}(u \otimes [u]^{\varepsilon}) \right) = 0. \quad (\text{A.11})$$

Next, notice from Lemma 2.3 that

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \operatorname{div}(u \otimes [u]^{\varepsilon}) = -\frac{1}{2} \int_0^T \int_{\Omega} \psi_{\tau} u \cdot \nabla \phi_{\delta} |[u]^{\varepsilon}|^2 \\ & \leq C \|u \cdot \nabla \phi_{\delta}\|_{L^2(0,T;L^2)} \int_0^T \int_{\{x: \operatorname{dist}(x, \partial\Omega) < \delta\}} |[u]^{\varepsilon}|^4 \rightarrow 0 \quad \text{as } \delta \rightarrow 0. \end{aligned} \quad (\text{A.12})$$

Therefore, (A.9) follows from (A.11) and (A.12). \square

Lemma A.2. *The pressure term in (A.8) satisfies*

$$\lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla P^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla \tilde{P}_i^{\varepsilon} \right) = 0. \quad (\text{A.13})$$

Proof. From (A.6), (A.7), and (A.10) we have

$$\|P\|_{L^2(\Omega)} \leq C. \quad (\text{A.14})$$

Next, we write

$$\begin{aligned} & \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla P^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla \tilde{P}_i^{\varepsilon} \right) \\ &= \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla P^{\varepsilon} + \sum_{i=1}^k \xi_i \nabla \tilde{P}_i^{\varepsilon} - \nabla [P]^{\varepsilon} \right) \\ & \quad + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} ([u]^{\varepsilon} - u) \cdot \nabla [P]^{\varepsilon} + \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} u \nabla [P]^{\varepsilon} \\ &=: K_1 + K_2 - \int_0^T \int_{\Omega} \psi_{\tau} u \cdot \nabla \phi_{\delta} [P]^{\varepsilon}. \end{aligned} \quad (\text{A.15})$$

The terms on the right-hand side will be treated as follows:

First, by (2.9), (A.2) and (A.14),

$$\begin{aligned} K_1 &= \int_0^T \int_{\Omega} \psi_{\tau} \phi_{\delta} [u]^{\varepsilon} \left(\xi_0 \nabla (P^{\varepsilon} - P) + \sum_{i=1}^k \xi_i \nabla (\tilde{P}_i^{\varepsilon} - P) \right) \\ &= - \int_0^T \int_{\Omega} \psi_{\tau} (\nabla \phi_{\delta} [u]^{\varepsilon} \xi_0 + \phi_{\delta} \operatorname{div} [u]^{\varepsilon} \xi_0 + \phi_{\delta} [u]^{\varepsilon} \nabla \xi_0) (P^{\varepsilon} - P) \\ & \quad - \sum_{i=1}^k \int_0^T \int_{\Omega} \psi_{\tau} (\nabla \phi_{\delta} [u]^{\varepsilon} \xi_i + \phi_{\delta} \operatorname{div} [u]^{\varepsilon} \xi_i + \phi_{\delta} [u]^{\varepsilon} \nabla \xi_i) (\tilde{P}_i^{\varepsilon} - P) \\ &\leq C(\delta) \left(\|P^{\varepsilon} - [P]^{\varepsilon}\|_{L^2_{loc}(0,T;L^2(V_0))} + \sum_{i=1}^k \|\tilde{P}_i^{\varepsilon} - [P]^{\varepsilon}\|_{L^2_{loc}(0,T;L^2(V_i))} \right) \end{aligned}$$

which implies $\lim_{\varepsilon \rightarrow 0} K_1 = 0$ due to (A.10) and (2.12).

The limit $\lim_{\varepsilon \rightarrow 0} K_2 = 0$ can be proved using similar arguments.

Finally, by (1.2), (A.2), the Hardy inequality, and (A.14) it follows that

$$\int_0^T \int_{\Omega} \psi_{\tau} u \cdot \nabla \phi_{\delta} P \leq C \|u \cdot \nabla \phi_{\delta}\|_{L^2(0,T;L^2)} \int_0^T \int_{\{x: \operatorname{dist}(x, \partial\Omega) < \delta\}} |P|^2 \rightarrow 0 \quad \text{as } \delta \rightarrow 0,$$

proving the lemma. \square

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