

Chem 2430 - Answers Exam #1.

1. Yes, e^{-ax^2} is reasonable as it $\rightarrow 0$ as $x \rightarrow \pm\infty$, and it is symmetric about $x=0$

$$\langle E \rangle = \frac{2 \int_0^\infty e^{-ax^2} H e^{-ax^2} dx}{2 \int_0^\infty e^{-2ax^2} dx}$$

$$\langle E \rangle = \frac{2 \int_0^\infty e^{-ax^2} \left(-\frac{1}{2} \frac{d^2}{dx^2} \right) e^{-ax^2} dx + 2V_0 \int_0^\infty \frac{1}{2} e^{-2ax^2} dx}{\int_0^\infty e^{-2ax^2} dx}$$

$$\langle E \rangle = \frac{2 \int_0^\infty \left(a - 2a^2 x^2 \right) e^{-2ax^2} dx + 2V_0 \int_0^\infty \frac{1}{2} e^{-2ax^2} dx}{\int_0^\infty e^{-2ax^2} dx}$$

$$= \frac{\frac{1}{2} \sqrt{a} \sqrt{\frac{\pi}{2}} + \sqrt{\frac{\pi}{2}} \frac{V_0}{\sqrt{a}} \operatorname{erfc} \left(\sqrt{\frac{a}{2}} b \right)}{\sqrt{\frac{\pi}{2}} \frac{1}{\sqrt{a}}}$$

$$= a/2 + V_0 \operatorname{erfc} \left(\sqrt{\frac{a}{2}} b \right)$$

A good approximate wavefunction for the first excited state is $x e^{-ax^2}$.

2. $H = \hbar\omega \left(a^\dagger a + \frac{1}{2} \right)$

$$[a, [1/2, a^\dagger]] = a \frac{1}{2} a^\dagger - a a^\dagger \frac{1}{2} - \frac{1}{2} a^\dagger a + \frac{1}{2} a^\dagger a = 0$$

$$[a, [a^\dagger a, a^\dagger]] = a a^\dagger a a^\dagger - a a^\dagger a^\dagger a - a^\dagger a a^\dagger a + a^\dagger a^\dagger a a$$

$$= a a^\dagger [a a^\dagger - a^\dagger a] - a^\dagger [a a^\dagger - a^\dagger a] a$$

$$\text{but } a a^\dagger - a^\dagger a = 1$$

$$\text{So } [a, [a^\dagger a, a^\dagger]] = a a^\dagger - a^\dagger a = 1$$

So $[a, [H, a^\dagger]] = \hbar\omega$ which is no longer an operator

3. I, $A e^{ik_1 x} + B e^{-ik_1 x}$, $k_1 = \sqrt{2mE}/\hbar$
 II, $C e^{ik_2 x} + D e^{-ik_2 x}$, $k_2 = \sqrt{2m(E-V_1)}/\hbar$
 III, $E e^{-Kx}$, $K = \sqrt{2m(V_2-E)}/\hbar$

4. In Cartesian coordinates

$$E = \frac{\hbar\omega}{2} + n_x \hbar\omega + \frac{\hbar\omega}{2} + n_y \hbar\omega = \hbar\omega (1 + n_x + n_y)$$

$$E(0,0) = \hbar\omega, \quad E(1,0) = E(0,1) = 2\hbar\omega$$

$$E(2,0) = E(0,2) = E(1,1) = 3\hbar\omega$$

In polar coordinates

$$-\frac{1}{2\mu} \left[\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right] \psi + \frac{1}{2} k r^2 \psi = E \psi$$

$$\psi = f(r) g(\phi) = f(r) e^{im\phi}$$

$$-\frac{1}{2\mu} \left[\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} - \frac{m^2}{r^2} \right] \psi + \frac{1}{2} k r^2 \psi = E \psi$$

Lets try $f(r) = e^{-ar^2}$, for the case $m=0$

This gives $E = 2a/\mu$, $a = \frac{1}{2} \sqrt{k\mu}$

$E = \sqrt{k/\mu} \rightarrow \hbar\omega$ same as in cartesian coord.

For $m = \pm 1$, try $f(r) = r e^{-ar^2}$

This gives $E = 2\hbar\omega$, again the same as above