

Operator approach to the Harmonic Oscillator

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{k\hat{x}^2}{2} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m\omega_0^2 \hat{x}^2$$

Change of variables

$$\hat{\zeta} = \beta \hat{x} \quad \text{where} \quad \beta = \sqrt{\frac{m\omega_0}{\hbar}} \quad \left| \quad \begin{array}{l} \text{dimensionless position} \\ \\ \text{dimensionless momentum} \end{array} \right.$$
$$\hat{\pi} = \frac{\beta}{m\omega_0} \hat{p}$$

$$\hat{H} = \frac{\hbar\omega_0}{2} [\hat{\pi}^2 + \hat{\zeta}^2]$$

$$\hat{h} = \frac{1}{\hbar\omega_0} \hat{H} = \frac{1}{2} [\hat{\pi}^2 + \hat{\zeta}^2]$$

$$\hat{H} |\psi\rangle = E |\psi\rangle \quad \Rightarrow \quad \hat{h} |\psi\rangle = \varepsilon |\psi\rangle$$

$$\varepsilon = \frac{E}{\hbar\omega_0} \quad \text{dimensionless energy}$$

$$[\hat{x}, \hat{p}] = i\hbar, \quad [\hat{\zeta}, \hat{\pi}] = i$$

$$[\hat{\zeta} + i\hat{\pi}][\hat{\zeta} - i\hat{\pi}] = \hat{\zeta}^2 + \hat{\pi}^2 - 1$$

$$\text{So } \hat{h} = \left(\frac{\hat{\zeta} - i\hat{\pi}}{\sqrt{2}} \right) \left(\frac{\hat{\zeta} + i\hat{\pi}}{\sqrt{2}} \right) + \frac{1}{2}$$

$$\text{raising operator } \hat{a}^\dagger = \left(\frac{\hat{\zeta} - i\hat{\pi}}{\sqrt{2}} \right)$$

$$\text{lowering operator } \hat{a} = \left(\frac{\hat{\zeta} + i\hat{\pi}}{\sqrt{2}} \right)$$

\hat{a} and \hat{a}^\dagger are Hermitian conjugates of one another

$$[\hat{a}, \hat{a}^\dagger] = 1$$

$$\hat{a}^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^\dagger + \hat{a})$$

$$\hat{p} = \sqrt{\frac{\hbar m\omega}{2}} i(\hat{a}^\dagger - \hat{a})$$

So $\hat{h} = \hat{a}^\dagger \hat{a} + \frac{1}{2}$

Can show that $\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle$

where $|n\rangle$ and $|n+1\rangle$ are the n^{th} and the $(n+1)^{\text{th}}$ eigenfunctions of \hat{h}

Also $\hat{a} |n\rangle = \sqrt{n} |n-1\rangle$

Note $\hat{a}^\dagger \hat{a} |n\rangle = \hat{a}^\dagger \sqrt{n} |n-1\rangle = n |n\rangle$

$\hat{N} = \hat{a}^\dagger \hat{a}$ Informs you of what level the system is in

So $\hat{h} |n\rangle = \left[(\hat{a}^\dagger \hat{a}) + \frac{1}{2} \right] |n\rangle = \left(n + \frac{1}{2} \right) |n\rangle$

Recall $\hat{H} = \hbar \omega_0 \hat{h}$

So the energy levels go as $(n + \frac{1}{2}) \hbar \omega_0$

Exactly what we found in solving the DE for the HO problem

Doing integrals

$$\langle n | \hat{\zeta} | n \rangle = ? \quad \left| \quad \hat{\zeta} = \frac{\hat{a}^\dagger + \hat{a}}{\sqrt{2}} \right.$$

$$\langle n | \frac{\hat{a}^\dagger + \hat{a}}{\sqrt{2}} | n \rangle = \frac{1}{\sqrt{2}} (\langle n | n+1 \rangle \sqrt{n+1} + \langle n | n-1 \rangle \sqrt{n}) = 0$$

$$\langle n | \zeta | m \rangle = \frac{1}{\sqrt{2}} [\sqrt{n} \delta_{m,n-1} + \sqrt{n+1} \delta_{m,n+1}]$$

$$\langle m | \hat{\zeta}^2 | m \rangle = ? \quad \left| \quad \begin{aligned} \hat{\zeta} &= \frac{\hat{a}^\dagger + \hat{a}}{\sqrt{2}} \\ \hat{\zeta}^2 &= \frac{\hat{a}^\dagger \hat{a}^\dagger + \hat{a} \hat{a} + \hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger}{2} \end{aligned} \right.$$

Or we can make use of $I = \sum |n\rangle\langle n|$

$$\langle m | \zeta^2 | m \rangle = \sum_n \langle m | \zeta | n \rangle \langle n | \zeta | m \rangle$$

$$a^\dagger |0\rangle = |1\rangle, a^\dagger |1\rangle = \sqrt{2} |2\rangle, a^\dagger |2\rangle = \sqrt{3} |3\rangle$$

$$|2\rangle = \frac{1}{\sqrt{2}} (\hat{a}^\dagger)^2 |0\rangle, |3\rangle = \frac{1}{\sqrt{6}} (\hat{a}^\dagger)^3 |0\rangle$$

$$\phi_n(\zeta) = \frac{(\hat{a}^\dagger)^n}{\sqrt{n!}} \phi_0(\zeta) = \frac{1}{\sqrt{n!}} \left(\frac{1}{\sqrt{2}} \zeta - \frac{1}{\sqrt{2}} \frac{d}{d\zeta} \right)^n \phi_0(\zeta)$$

Can show that $\phi_0(\zeta) = A_0 e^{-\zeta^2/2}$

and eventually get to
$$\phi_n(\zeta) = \frac{1}{n! 2^n \sqrt{\pi}} e^{-\zeta^2/2} H_n(\zeta)$$

(different normalization constant if using variable x)

Often one finds the step-up and step-down operators defined directly in terms of x and y .

$$\hat{a}^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^\dagger + \hat{a})$$

$$\hat{p} = \sqrt{\frac{\hbar m\omega}{2}} i (\hat{a}^\dagger - \hat{a})$$