CHAPTER 24 INTERMOLECULAR INTERACTIONS

potential energy u(r) between two particles

$$force = f = -\frac{du}{dr}$$

 $r = r^* =$ potential energy minimum

long-range coulombic

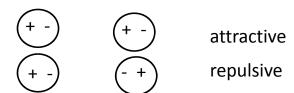
$$u \sim \frac{q_1 q_2}{r}$$
, q_1 , q_2 = charges on the two species

short-ranged van der Waals

$$u \sim -\frac{C}{r^6}$$

At short distance, u must repel one another $\left(Ae^{-Br};Ar^{-12}\right)$

dipole-dipole $\sim \frac{1}{r^3}$ long-range interaction between two H₂O molecules



charge-polarization ~ $-\frac{q\alpha}{2r^4}$ (e.g., Na⁺ Ar), α = polarizability of neutral species

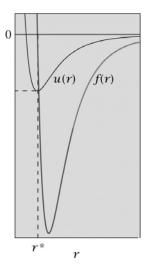


Figure 24.1 Molecular Driving Forces 2/e (© Garland Science 2011

u(r) 0 r^{-6} $-r^{-1}$

Figure 24.2 Molecular Driving Forces 2/e (© Garland Science 2011)

permanent moments

dipole (e.g., HF) $(\mu = qR)$ quadrupole (e.g., CO₂) octapole (e.g., CH₄)

charge-dipole
$$u \sim \frac{1}{r^2}$$

average over orientations $u \sim \frac{1}{3kT} \frac{1}{r^4}$

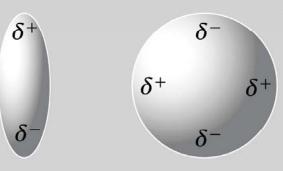
dipole-dipole

average over orientations $u \sim \frac{1}{3kT} \frac{1}{r^6}$

(a) Dipole

(b) Quadrupole





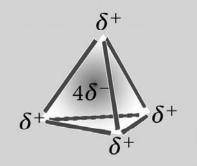


Figure 24.3 Molecular Driving Forces 2/e (© Garland Science 2011)

How do we understand the attraction between two Ar atoms?

London-dispersion

$$u = \frac{C_6}{r^6}$$
: $C_6 \propto \alpha_A \alpha_B$

 $\alpha_A, \alpha_B =$ polarizabilities of the two atoms/molecules

Lennard Jones potential

$$u(r) = \frac{a}{r^{12}} - \frac{b}{r^6}$$

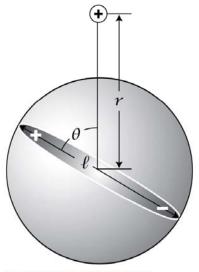


Figure 24.5 Molecular Driving Forces 2/e (© Garland Science 2011

H-bonds – have electrostatic, dispersion, and polarization contributions

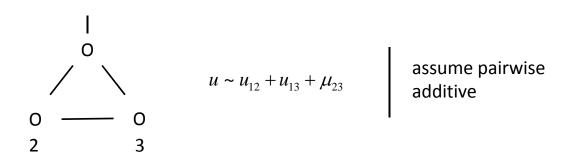








Figure 24.6 Molecular Driving Forces 2/e (© Garland Science 201

$$p = \frac{NkT}{V - Nb} - \frac{aN^2}{V^2} = \frac{\rho RT}{1 - b\rho} - a\rho^2$$

$$p = -\left(\frac{\partial F}{\partial V}\right)_{T,N} = -\left(\frac{\partial U}{\partial V}\right)_{T,N} + \left(\frac{\partial S}{\partial V}\right)_{T,N}$$

$$\frac{-aN^2}{V^2} \qquad \frac{NkT}{V - Nb} \qquad \text{from lattice model (Ex. 6.1)}$$
from attractions

particles in shell of radius r =

(density) x (volume of shell)

$$\rho$$
 x $4\pi r^2 dr$

$$U = \frac{N}{2} \int_0^\infty u(r) \rho 4\pi r^2 dr$$

assuming distribution is uniform

take $u(r) = \begin{cases} \infty, & r < r^* \\ -u_0 \left(\frac{r^*}{r}\right)^6 & r > r^* \end{cases}$

 $\int_0^\infty \to \int_{r^*}^\infty$ since no particles are between ϕ and r^*

$$U = -\frac{aN^2}{V}$$
, where $a = \frac{2\pi (r^*)^3}{3} u_0$

$$p = -\frac{aN^2}{V^2} - \frac{kT}{b_0} \ell n \left(1 - \frac{Nb_0}{V} \right)$$
, where $b_0 = \frac{V}{M} = \text{volume/# lattice sites}$

$$\approx \frac{NkT}{V - Nb} - \frac{aN^2}{V^2}$$
, $b = \frac{b_0}{2} = \frac{1}{2}$ volume of a particle

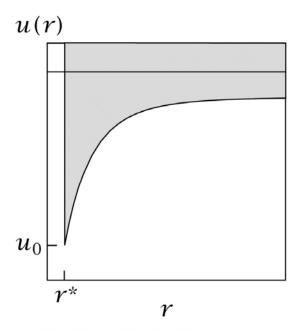


Figure 24.9 Molecular Driving Forces 2/e (© Garland Science 2011)

any potential falling off more rapidly than $\frac{1}{r^3}$ would give this result

Radial distribution functions

true density in shell at r from test particle

$$= \rho g(r)$$
 radial distribution function

where ρ = average density = N/V

$$g(r) = \rho^{true} / \rho^{avg}$$

molecules in 1st solvation shell =

$$\int_0^B \rho q(r) 4\pi r^2 dr$$

$$\int_0^\infty \rho q(r) 4\pi r^2 dr = n - 1 \approx n$$

$$U' = \int_0^B \rho u(r)g(r)4\pi r^2 dr =$$
 energy of interaction between one particle and all others

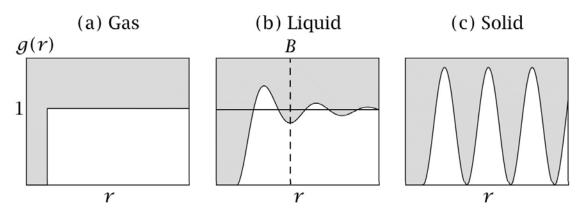


Figure 24.10 Molecular Driving Forces 2/e (© Garland Science 2011)

Lattice model

w (from Chapter 14) = $u(r^*)$

$$U = \frac{N}{2} \sum_{r=0}^{\infty} u(r) g(r) \rho 4\pi r^{2} = \frac{N}{2} u(r^{*}) z = \frac{Nwz}{2}$$

sum over nearest neighbors only

For molecules for which van der Waals (dispersion) interactions dominate:

$$w_{AA} = -c\alpha_A^2$$
 $w_{BB} = -cw_B^2$ $w_{AB} = -c\alpha_A\alpha_B$

c = constant ~ equal for molecules of ~ same size

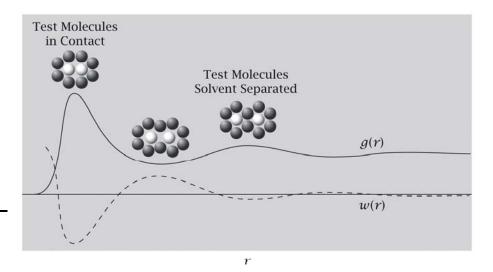


Figure 24.11 Molecular Driving Forces 2/e (© Garland Science 2011)