

$$\left\{ \begin{array}{l} \frac{1}{\Theta} \sin \theta \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \beta \sin^2 \theta = m_\ell^2 \\ \frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -m_\ell^2 \end{array} \right.$$

$$\Phi_{m_\ell} = A e^{im_\ell \phi}, \quad m_\ell = 0, \pm 1, \pm 2, \dots \quad (\text{but see below})$$

$$\left. \begin{array}{l} \beta = \ell(\ell + 1), \quad \ell = 0, 1, 2, \dots \\ m_\ell = -\ell, -\ell + 1, \dots, 0, \dots, \ell - 1, \ell \end{array} \right\} \text{quantization conditions}$$

$$\ell = 0 \quad \rightarrow \quad m_\ell = 0$$

$$\ell = 1 \quad \rightarrow \quad m_\ell = -1, 0, 1$$

$$\ell = 2 \quad \rightarrow \quad m_\ell = -2, -1, 0, 1, 2$$

s }
p } H atom
d }

$$Y(\theta, \phi) = Y_\ell^{m_\ell}(\theta, \phi) = \Theta_\ell^{m_\ell}(\theta) \Phi_{m_\ell}(\phi)$$

two boundary conditions \rightarrow two quantum #s

$$\beta = \frac{2\mu r_o^2 E}{\hbar^2} = \frac{2I}{\hbar^2} E$$

$$E = \frac{\hbar^2}{2I} \ell(\ell+1), \quad \ell = 0, 1, 2, \dots$$

$$\hat{H}Y_\ell^{m_\ell} = \frac{\hbar^2}{2I} \ell(\ell+1)Y_\ell^{m_\ell}$$

\uparrow degeneracy = $2\ell+1$

$$\hat{\ell}^2 Y_\ell^{m_\ell} = \hbar^2 \ell(\ell+1)Y_\ell^{m_\ell}$$

$\hat{\ell}^2$ and \hat{H} obviously commute

$\hat{\ell}^2$: “angular momentum”
operator

$$|\vec{\ell}| = \hbar \sqrt{\ell(\ell+1)}$$

components of the angular momentum operator

$$\hat{l}_x = \frac{\hbar}{i} \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) = \frac{\hbar}{i} \left(-\sin \phi \frac{\partial}{\partial \theta} - \cot \theta \cos \phi \frac{\partial}{\partial \phi} \right)$$

$$\hat{l}_y = \frac{\hbar}{i} \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) = \frac{\hbar}{i} \left(\cos \phi \frac{\partial}{\partial \theta} - \cot \theta \sin \phi \frac{\partial}{\partial \phi} \right)$$

$$\hat{l}_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

$$\left[\hat{l}_x, \hat{l}_y \right] = i\hbar \hat{l}_z$$

$$\left[\hat{l}_y, \hat{l}_z \right] = i\hbar \hat{l}_x$$

$$\left[\hat{l}_z, \hat{l}_x \right] = i\hbar \hat{l}_y$$

$$\hat{l}_z Y_\ell^{m\ell} = m_\ell \hbar Y_\ell^{m\ell}$$

Can simultaneously know the magnitude of the angular momentum and one of its components

Spherical harmonics

$$Y_0^0 = \frac{1}{\sqrt{4\pi}} \quad \text{spherically symmetric} \longrightarrow \quad s$$

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta \quad \longrightarrow \quad p_z$$

$$Y_1^{\pm 1} = \left(\frac{3}{8\pi}\right)^{1/2} \sin \theta e^{\pm i\phi} \quad \left\{ \begin{array}{l} \sqrt{\frac{3}{4\pi}} \sin \theta \cos \phi \quad \longrightarrow \quad p_x \\ \sqrt{\frac{3}{4\pi}} \sin \theta \sin \phi \quad \longrightarrow \quad p_y \end{array} \right.$$

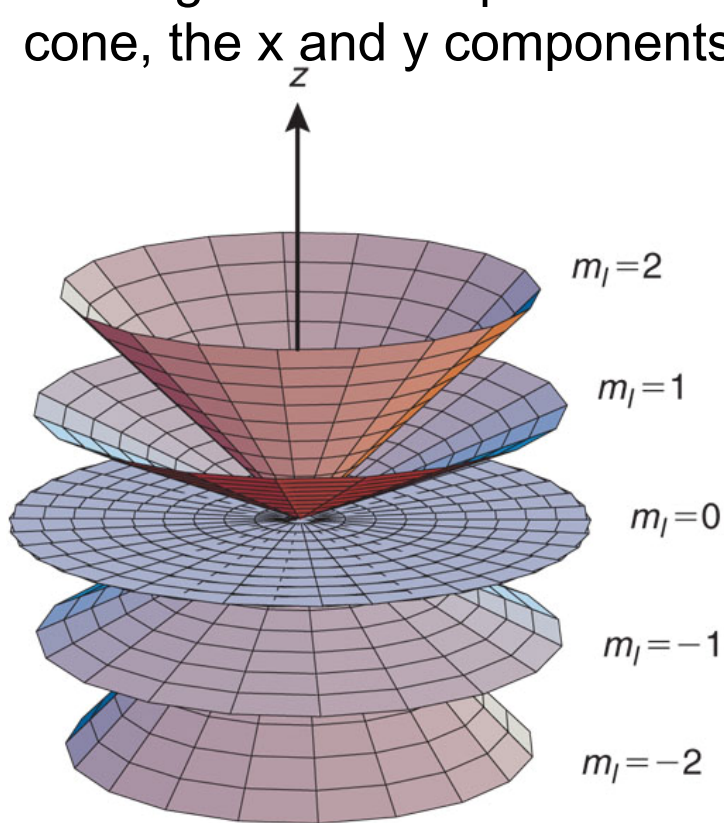
$$Y_2^0 = \left(\frac{5}{16\pi}\right)^{1/2} (3\cos^2 \theta - 1) \quad \longrightarrow \quad d_z^2$$

$$Y_2^{\pm 1} = \left(\frac{15}{8\pi}\right)^{1/2} \sin \theta \cos \theta e^{\pm i\phi} \quad \left\{ \begin{array}{l} \sqrt{\frac{15}{4\pi}} \sin \theta \cos \theta \cos \phi \quad \longrightarrow \quad d_{xz} \\ \sqrt{\frac{15}{4\pi}} \sin \theta \cos \theta \sin \phi \quad \longrightarrow \quad d_{yz} \end{array} \right.$$

$$Y_2^{\pm 2} = \left(\frac{15}{32\pi}\right)^{1/2} \sin^2 \theta e^{\pm 2i\phi} \quad \left\{ \begin{array}{l} \sqrt{\frac{15}{16\pi}} \sin^2 \theta \cos 2\phi \quad \longrightarrow \quad d_{x^2-y^2} \\ \sqrt{\frac{15}{16\pi}} \sin^2 \theta \sin 2\phi \quad \longrightarrow \quad d_{xy} \end{array} \right.$$

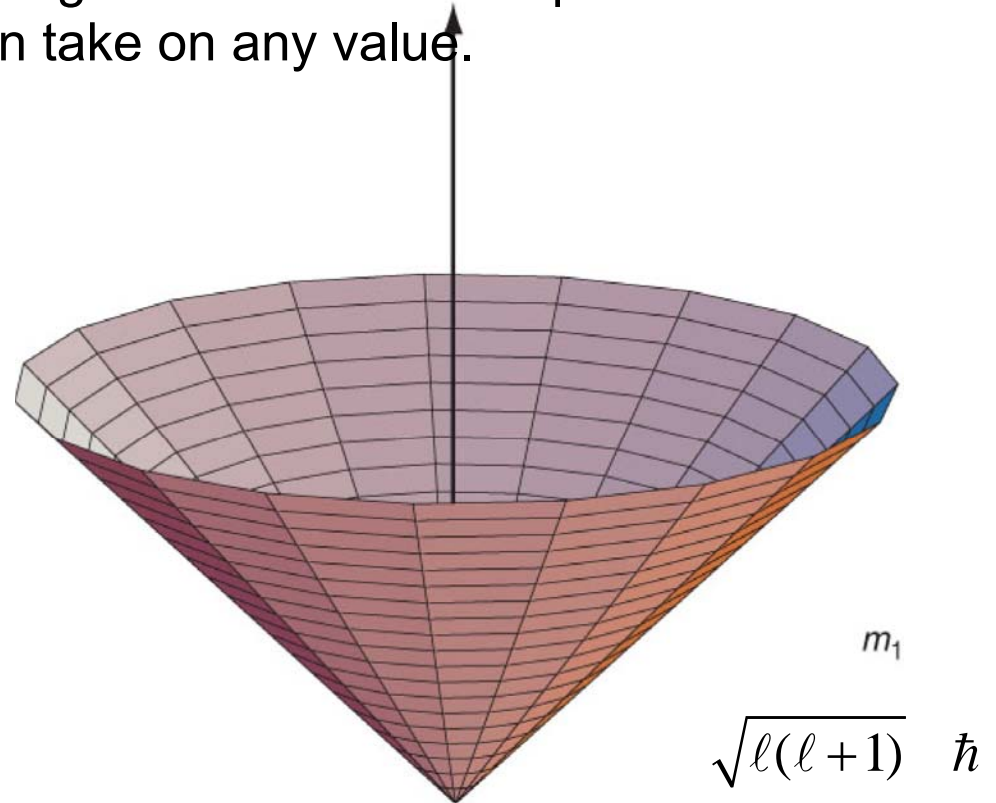
Consider a sphere of radius $\sqrt{\ell(\ell+1)}\hbar$

For $\ell=2$, the allowed solutions can be represented as 4 cones and 1 disk. Although the z component of the angular momentum is specified for each cone, the x and y components can take on any value.



$$\ell = 2$$

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