

## Chapter 4, continued

$$\left. \begin{aligned} \text{particle in 3D box, } V = 0 \text{ for } 0 < x < a, \\ 0 < y < b, \\ 0 < z < c \\ = \infty, \text{ otherwise} \end{aligned} \right\}$$

$$\frac{-\hbar^2}{2m} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi(x, y, z) = E\psi(x, y, z)$$

assume the problem separates

$$\psi(x, y, z) = X(x)Y(y)Z(z)$$

$$\frac{-\hbar^2}{2m} \left[ YZ \frac{d^2 X}{dx^2} + XZ \frac{d^2 Y}{dy^2} + XY \frac{d^2 Z}{dz^2} \right] = EXYZ$$

$$\frac{-\hbar^2}{2m} \left[ \underbrace{\frac{1}{X} \frac{d^2 X}{dx^2}} + \underbrace{\frac{1}{Y} \frac{d^2 Y}{dy^2}} + \underbrace{\frac{1}{Z} \frac{d^2 Z}{dz^2}} \right] = E$$

depends only on x    depends only on y    depends only on z

a constant

$$\Rightarrow \left. \begin{aligned} \frac{-\hbar^2}{2m} \frac{d^2 X}{dx^2} &= E_x X \\ \frac{-\hbar^2}{2m} \frac{d^2 Y}{dy^2} &= E_y Y \\ \frac{-\hbar^2}{2m} \frac{d^2 Z}{dz^2} &= E_z Z \end{aligned} \right\} E = E_x + E_y + E_z$$

$$E = \frac{h^2}{8m} \left( \frac{n_x^2}{a^2} + \frac{n_y^2}{b^2} + \frac{n_z^2}{c^2} \right), \quad n_x, n_y, n_z = 1, 2, \dots$$

$$\psi(x, y, z) = N \sin\left(\frac{n_x \pi x}{a}\right) \sin\left(\frac{n_y \pi y}{b}\right) \sin\left(\frac{n_z \pi z}{c}\right)$$

let  $a = b = c$

$$E = \frac{h^2}{8ma^2} (n_x^2 + n_y^2 + n_z^2)$$

$$E(1,1,1) = \frac{h^2}{8ma^2} \quad (3)$$

$$\begin{cases} E(2,1,1) = \frac{h^2}{8ma^2} \quad (6) \\ E(1,2,1) \\ E(1,1,2) \end{cases}$$

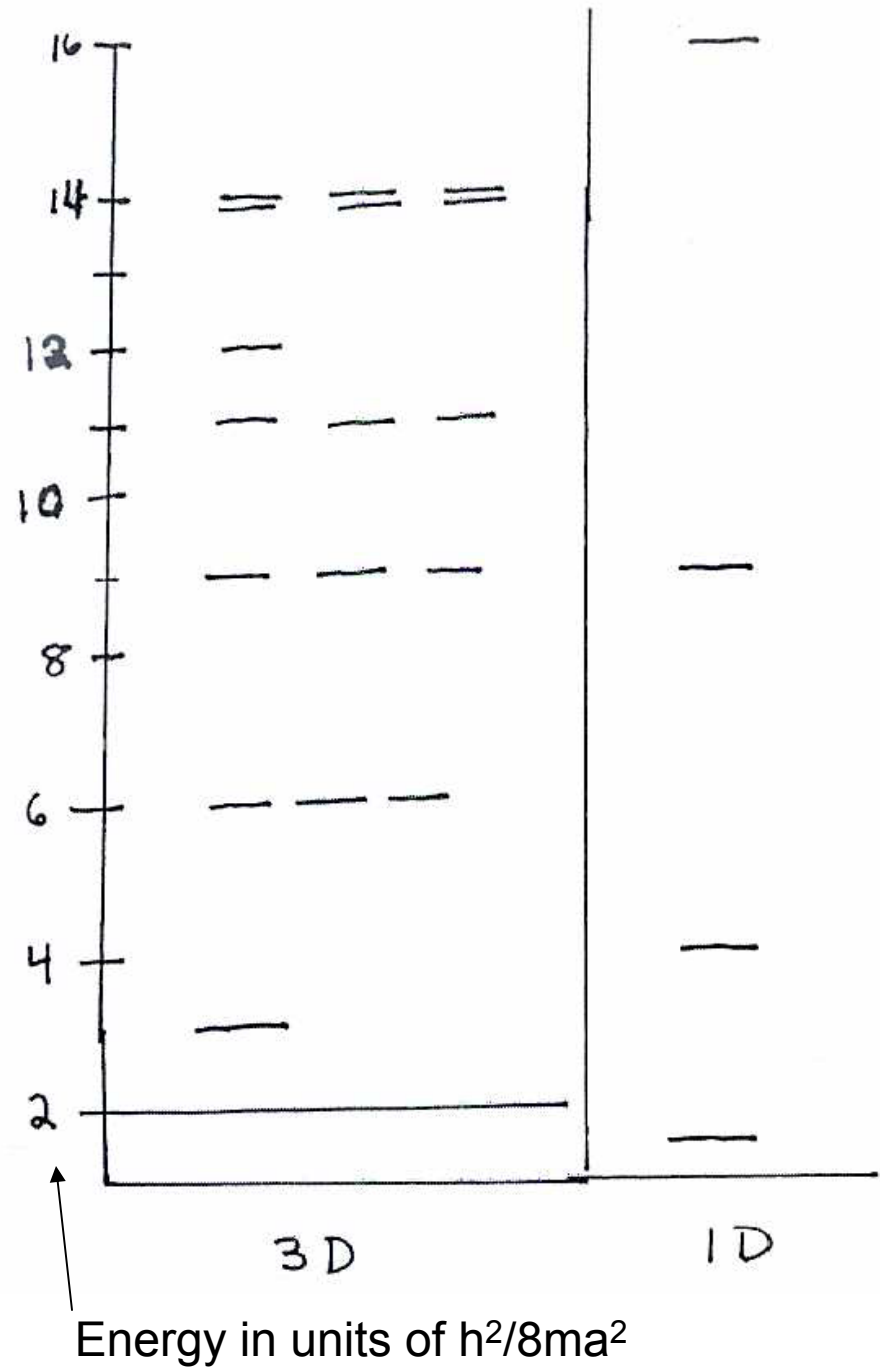
$$\begin{cases} E(2,2,1) = \frac{h^2}{8ma^2} \quad (9) \\ E(2,1,2) \\ E(1,2,2) \end{cases}$$

$$\begin{cases} E(3,1,1) = \frac{h^2}{8ma^2} \quad (11) \\ E(1,3,1) \\ E(1,1,3) \end{cases}$$

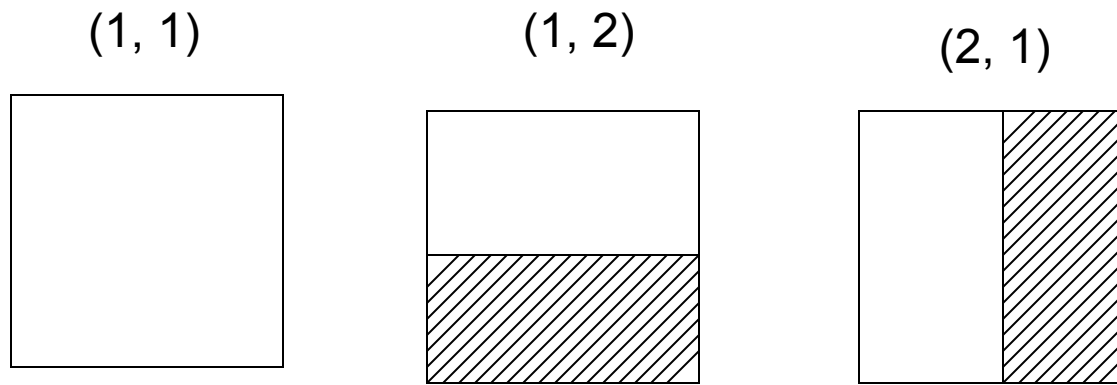
$$E(2,2,2) = \frac{h^2}{8ma^2} \quad (12)$$

$$E(1,2,3) = \frac{h^2}{8ma^2} \quad (14)$$

6-fold degenerate



- Note how much more rapidly energy levels grow for 3D vs. 1D
- Degeneracies are a result of symmetry



Nodal patterns for  $(1,1)$ ,  $(1,2)$ ,  $(2,1)$   
eigenfunctions of 2D particle-in-box  
problem

back to the 1D particle-in-box problem:

example 4.2

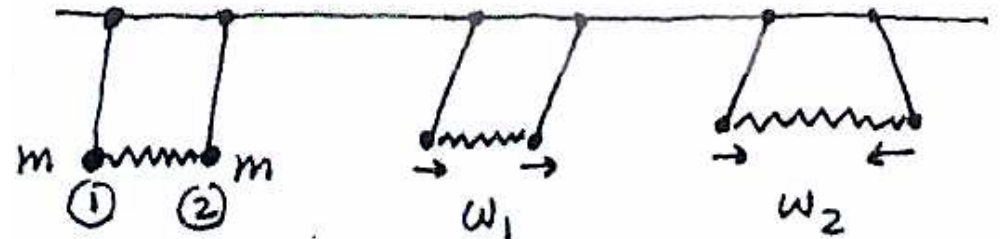
$$\psi = c \sin\left(\frac{\pi x}{a}\right) + d \sin\left(\frac{2\pi x}{a}\right) \longleftarrow \text{Not an e.f. of H unless } c \text{ or } d = 0$$

$$\Psi(x,t) = c e^{-iE_1 t/\hbar} \sin\left(\frac{\pi x}{a}\right) + d e^{-iE_2 t/\hbar} \sin\left(\frac{2\pi x}{a}\right)$$

$$\neq \psi(x) f(t)$$

$\Rightarrow$  Not a standing wave

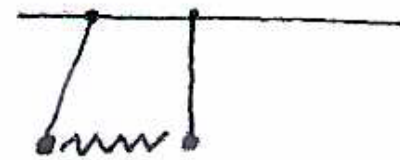
### Classical analog



two fundamental frequencies

$\omega_1, \omega_2$

what happens with the initial condition shown to the right (1 is displaced but 2 is not)?

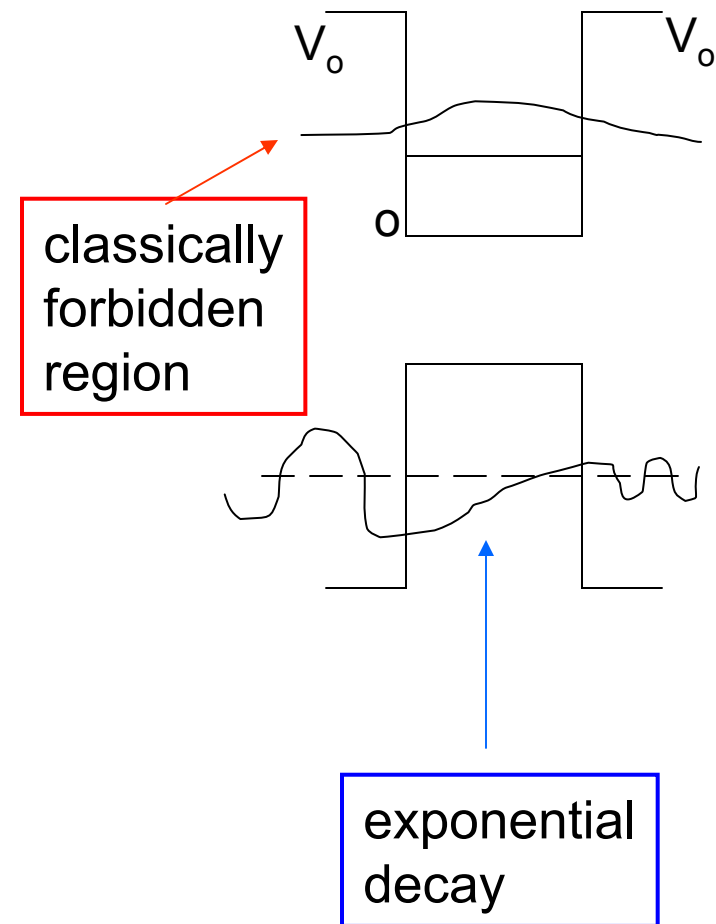


What happens if the box is finite?

The wavefunction now leaks (tunnels) outside the box

What if there is a barrier?

The particle can tunnel through the barrier



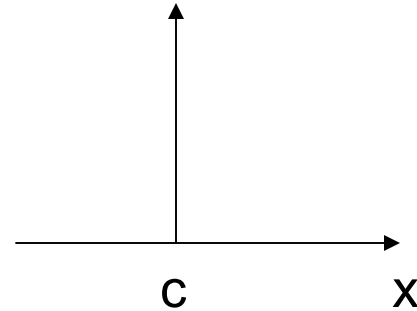
### Some additional exercises:

$$\text{if } \psi(x) = \sqrt{\frac{2}{a}} \left[ c \sin \frac{\pi x}{a} + d \sin \frac{2\pi x}{a} \right], \quad c^2 + d^2 = 1$$

$$\text{What is } \langle \hat{H} \rangle? \quad c^2 E_1 + d^2 E_2 = c^2 \frac{\hbar^2}{8ma^2} + d^2 \frac{4\hbar^2}{8ma^2}$$

$$\begin{aligned} \langle \psi | \hat{H} | \psi \rangle &= \frac{2}{a} \int_0^a \left( c \sin \frac{\pi x}{a} + d \sin \frac{2\pi x}{a} \right) \left( \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right) \left( c \sin \frac{\pi x}{a} + d \sin \frac{2\pi x}{a} \right) \\ &= \frac{2}{a} \left[ \int_0^a \left( c \sin \frac{\pi x}{a} + d \sin \frac{2\pi x}{a} \right) \left( \frac{\hbar^2}{2m} \frac{\pi^2}{a^2} \right) \left( c \sin \frac{\pi x}{a} + 4d \sin \frac{2\pi x}{a} \right) \right] \\ &= \frac{\hbar^2}{4ma^3} \left[ c^2 \int_0^a \sin^2 \frac{\pi x}{a} dx + 4d^2 \int_0^a \sin^2 \frac{2\pi x}{a} dx \right] \\ &= \frac{\hbar^2}{8ma^2} [c^2 + 4d^2] \end{aligned}$$

$$\begin{aligned}\delta(x-c) &= \text{delta function} = 0, \quad x \neq c \\ &= \infty, \quad x = c\end{aligned}$$



$$\int_{-\infty}^{\infty} f(x)\delta(x-c)dx = f(c)$$

Convince yourself that for the particle-in-box problem,  $\delta(x-c)$  can be represented as a sum over all  $\sin\left(\frac{n\pi x}{a}\right)$  functions.

$\Rightarrow$  momentum ranges over all possible values

Note the connection with the uncertainty principle