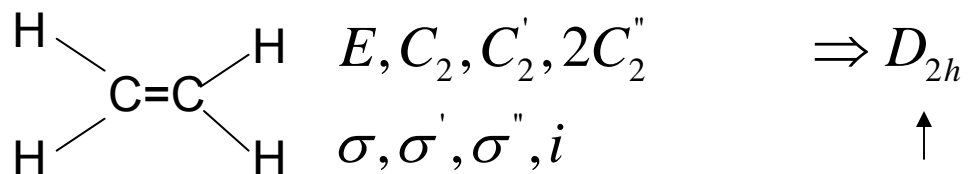
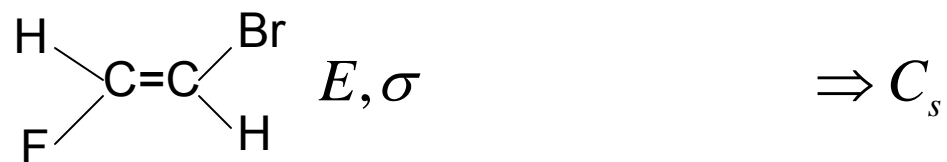


Chapter 17 – Symmetry

Symmetry elements

E	-	identity
C_n	-	n-fold rotation
σ	-	mirror plane
i	-	inversion
S_n	-	n-fold rotation-reflection

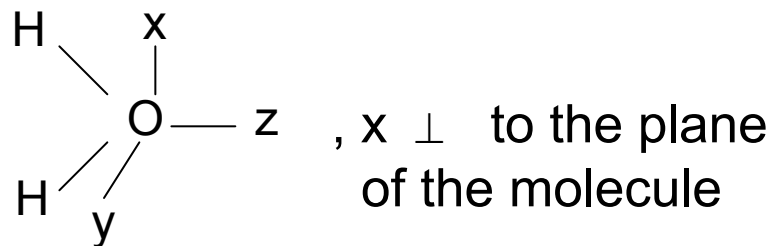
H_2O	$E, \sigma_v, \sigma_v', C_2$	$\Rightarrow C_{2v}$	$\sigma_v \Rightarrow$ mirror plane contains princ. rotational axis
NH_3	$E, \sigma_v(1), \sigma_v(2), \sigma_v(3)$ $C_3, C_3^2 = C_3^{-1}$	$\Rightarrow C_{3v}$	$\sigma_h \Rightarrow$ mirror plane \perp to the princ. axis
$CHClBr$	E	$\Rightarrow C_1$	



↑
group

Symmetry operators can be represented by matrices

Example:



$$\hat{E}: \begin{pmatrix} 1 & 0 & 1 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{C}_2: \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{\sigma}_v: \begin{pmatrix} 1 & 0 & 1 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{\sigma}'_v: \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{C}_2 \cdot \hat{C}_2 = \hat{E}$$

$$\hat{C}_2 \cdot \hat{\sigma}_v = \hat{\sigma}'_v$$

etc.

TABLE 17.3

Multiplication
Table for Operators
of the C_{2v} Group

Second Operation	First Operation			
	\hat{E}	\hat{C}_2	$\hat{\sigma}_v$	$\hat{\sigma}'_v$
\hat{E}	\hat{E}	\hat{C}_2	$\hat{\sigma}_v$	$\hat{\sigma}'_v$
\hat{C}_2	\hat{C}_2	\hat{E}	$\hat{\sigma}'_v$	$\hat{\sigma}_v$
$\hat{\sigma}_v$	$\hat{\sigma}_v$	$\hat{\sigma}'_v$	\hat{E}	\hat{C}_2
$\hat{\sigma}'_v$	$\hat{\sigma}'_v$	$\hat{\sigma}_v$	\hat{C}_2	\hat{E}

form a representation
from the C_{2v} group

In this case, there are simpler representations

C_{2v}	\hat{E}	\hat{C}_2	$\hat{\sigma}_2$	$\hat{\sigma}'_v$		
A_1	1	1	1	1	z	x^2, y^2, z^2
A_2	1	1	-1	-1	R_z	xy
B_1	1	-1	1	-1	x, R_x	xz
B_2	1	-1	-1	1	y, R_x	yz

} character table



irreducible representations

$$O \quad p_z \rightarrow a_1$$

A_1 = totally symmetric representation

$$O \quad p_x \rightarrow b_1$$

$$O \quad p_y \rightarrow b_2$$

for C_{2v} all irreducible representations are one-dimensional

\Rightarrow no degeneracies

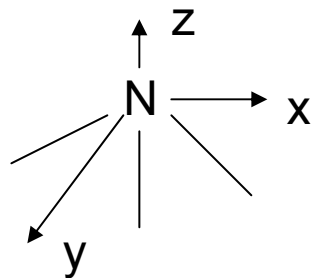
C_{3V} is an example of a group with a degenerate representation

TABLE 17.4

The C_{3V} Character Table

	E	$2C_3$	$3\sigma_v$		
A_1	1	1	1	z	$x^2 + y^2, z^2$
A_2	1	1	-1	R_z	
E	2	-1	0	$(x, y), (R_x, R_y)$	$(x^2 - y^2, xy), (xz, yz)$

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rotate 120° , "mixes" x and y

E is a two-fold degenerate representation

The different representations are orthogonal

$$A_1 \times A_2 = 1 \cdot 1 + 2 \cdot 1 \cdot 1 + 3 \cdot 1 \cdot (-1) = 0$$

$$A_2 \times E = 1 \cdot 2 + 2(1)(-1) + 0 = 0$$

Electronic structure

C_{2v} group

$$a_2^2 \rightarrow A_1$$

$$b_1 b_2 \rightarrow A_2$$

$$a_2 b_2 \rightarrow A_1$$

$$b_1 a_2 \rightarrow B_2$$

$$\int \psi_1 \hat{H} \psi_2 d\tau = 0$$

if ψ_1, ψ_2 not the
same symmetry

$$\int \psi_1 \hat{A} \psi_2 d\tau = 0$$

if $\psi_1 \hat{A} \psi_2$ does not
contain totally symmetric
representation

Selection rules $\int a_1 z a_1 d\tau \neq 0$

$$a_1 \rightarrow a_1$$

$$\int a_1 z b_2 d\tau = 0$$

$$a_1 \rightarrow b_1$$

$$\int a_1 x b_1 d\tau \neq 0$$

etc.
allowed
transitions

$$\begin{aligned}
C_{3v} \quad e^2 &\rightarrow 4 \quad 1 \quad 0 \\
&= C_1(1 \quad 1 \quad 1) + C_2(1 \quad 1 \quad -1) + C_3(2 \quad -1 \quad 0) \\
C_1 = C_2 = C_3 = 1 &\Rightarrow e^2 \rightarrow e, a_1, a_2
\end{aligned}$$

two electrons in an e orbital $\rightarrow E, A_1, A_2$ states

Symmetries of vibrational normal modes

$$V = \frac{1}{2} \sum_i \left(\frac{\partial^2 V}{2Q_i^2} \right) Q_i^2, \quad Q_i \text{ are normal coordinates}$$

$$\Psi = \psi_1(Q_1) \psi_2(Q_2) \dots \psi_N(Q_N)$$

$$E = \sum_j \left(n_j + \frac{1}{2} \right) h\nu_j$$

Procedure (p. 404-405) for determining how many normal modes there are of each symmetry

$$H_2O = A_1, A_1, B_2$$

$$\psi_0(Q_j) Q_j \underbrace{\psi_m(Q_j)}$$

must belong to the same representation as x, y, or z to be IR active

C_{2v} : A_1 vibrations are IR active
 B_2 vibration is also IR active

T_d : example CH_4 : $A_1, E, 2T_2$ vibrations

IR
forbidden

IR
active